

Consistency of Ice Cloud Models in Forward Retrieval and Radiative Forcing Assessment

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Contributions to this presentation: M. Saito, S. Hioki, and C. P. Kuo

Cirrus (Ci)



Cirrostratus (Cs)



Cirrocumulus (Cu)



Altostratus (As)



Altostratus (As)



Stratocumulus (Sc)



Stratus (St)

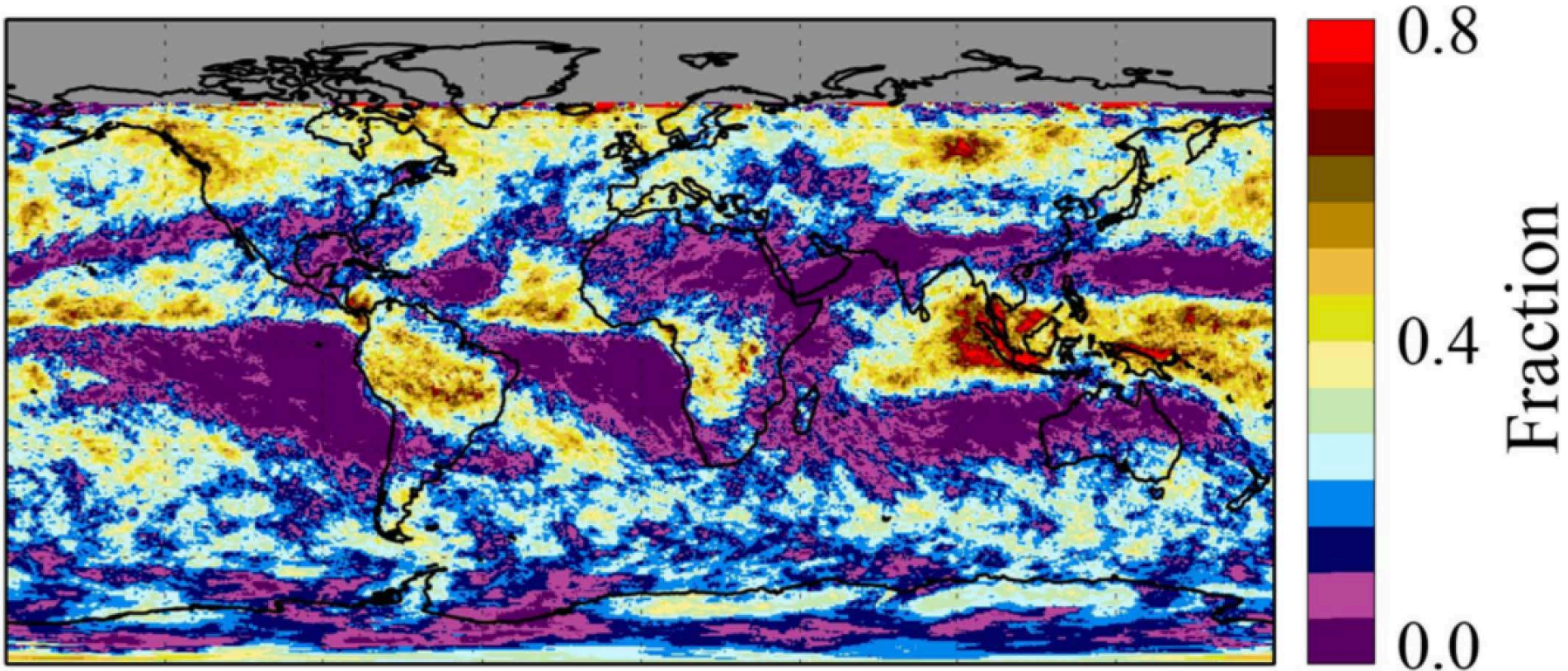


Cumulus (Cu, Fairweather)



Taken from <http://www.clouds-online.com>. Copyright information: <http://www.clouds-online.com/imprint.htm> except for Cirrus (taken from http://www.c-f-r.dk/images/Artikelbilleder_540_200px/cirrus_over_warsaw_june_26_2005.jpg) & Altostratus (<http://cimss.ssec.wisc.edu/satmet/gallery/images/altostratus.jpg>).

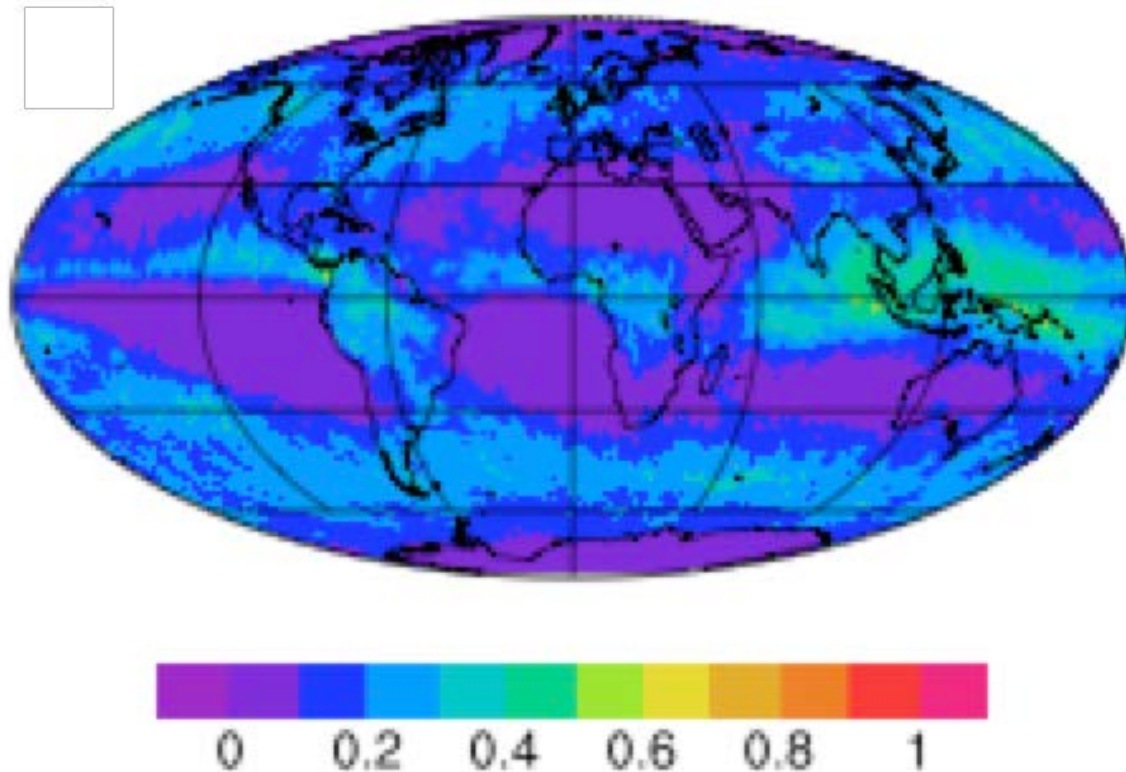
Ice Cloud Fraction based on MODIS Collection 6 data



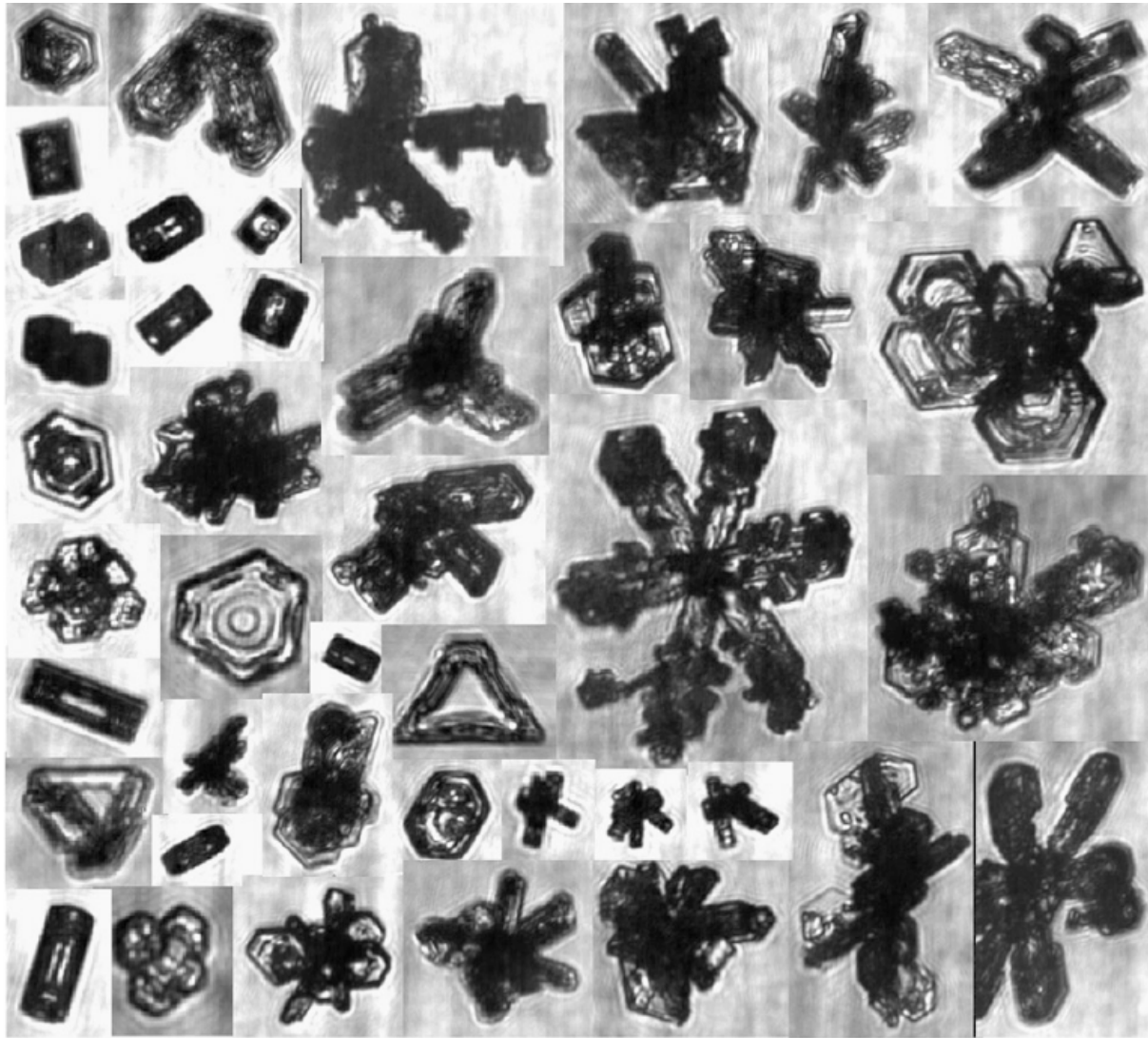
November 2012 Aqua MODIS monthly cloud fraction
([Platnick et al. 2017](#))

Global Ice Cloud Coverage

CF C6, Glb avg: 0.17



Ice cloud coverage (Yi et al. 2017) based on **one year (2012)** of level-2 MODIS Collection 6 cloud products (Platnick, Meyer et al. 2016)

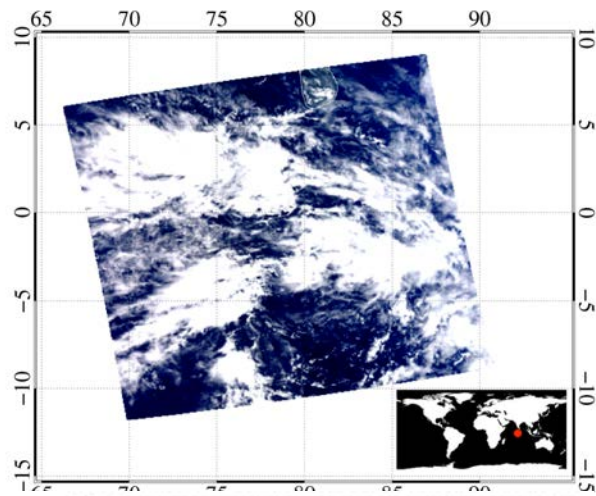


Bailey, M. P., and J. Hallett, 2009: A comprehensive habit diagram for atmospheric ice crystals: Confirmation from the laboratory, AIRS II, and other field studies, *J. Atmos. Sci.*, 66(9), 2888-2899, doi=10.1175/2009JAS2883.1.

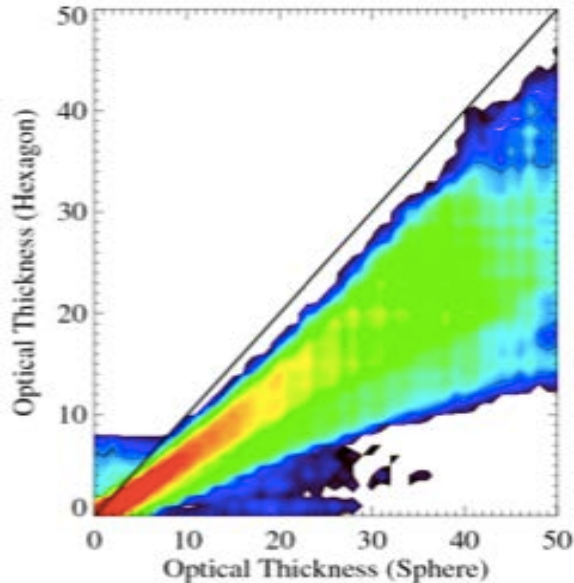
Stephens, G. L., S.-C. Tsay, P. W. Stackhouse Jr., and P. J. Flatau, 1990: The relevance of the microphysical and radiative properties of cirrus clouds to climate and climatic feedback. J. Atmos. Sci., 47, 1742-1754.

- ★ “The asymmetry parameter had to be **adjusted from** the broadband **Mie value** of **$g=0.87$** for the size distribution chosen to a lower value of **$g=0.7$** in order to **bring the observations and theory into broad agreement.**”
- ★ “Cirrus clouds characterized by **$g=0.87$** warmed **approximately twice** as much as cirrus clouds modeled **with $g=0.7$.**”

Retrieved optical thickness based on MODIS observations

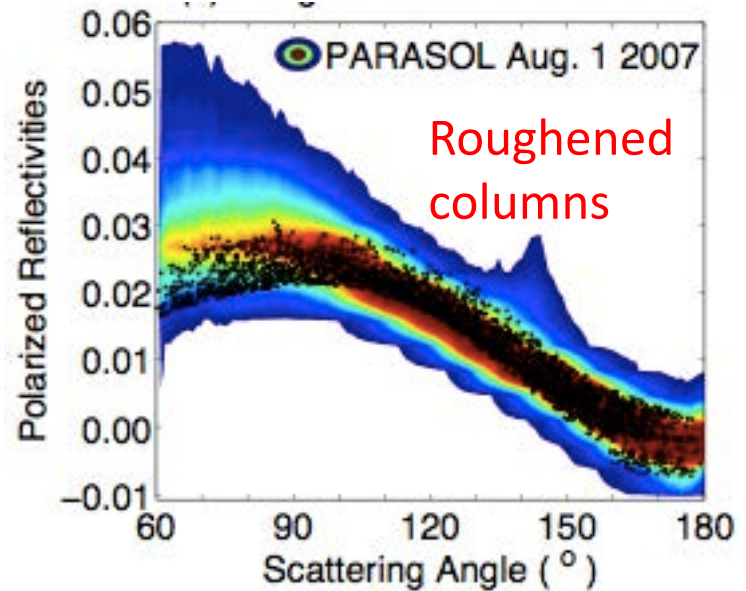
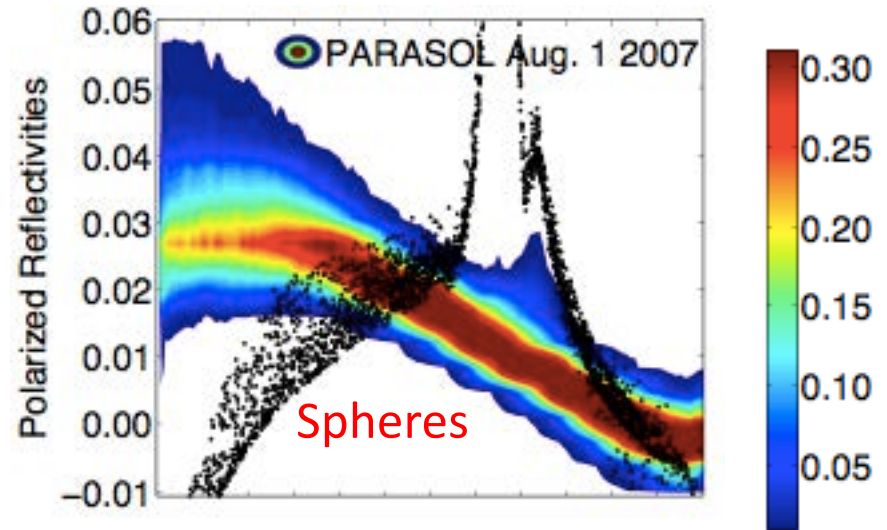


Hexagonal Ice Crystals



Ice Spheres

Simulated polarized reflectivity vs POLDER observations



Technical readiness

*Modeling capabilities for
computing the optical
properties of nonspherical ice
crystals*

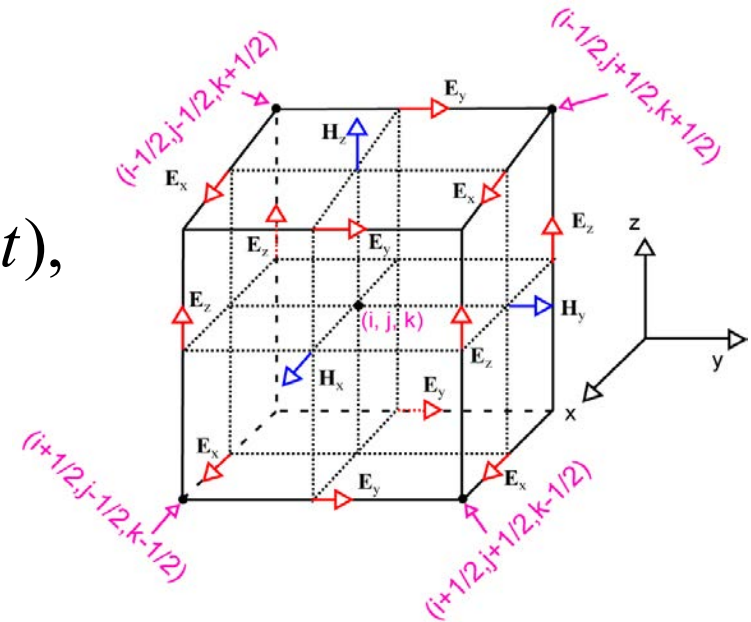
Finite-difference Time Domain (FDTD) Method

(Yee 1966; Taflove and Hagness 2000; Yang and Liou, 1996; ...)

Second order central difference scheme applied to the time-dependent Maxwell curl equations:

$$\nabla \times \mathbf{H}(\mathbf{r}, t) = \frac{\varepsilon}{c} \frac{\partial \mathbf{E}(\mathbf{r}, t)}{\partial t} + \frac{4\pi}{c} \sigma \mathbf{E}(\mathbf{r}, t),$$

$$\nabla \times \mathbf{E}(\mathbf{r}, t) = -\frac{1}{c} \frac{\partial \mathbf{H}(\mathbf{r}, t)}{\partial t},$$



Staggered locations of the E & H field components on a cubic cell (Yee, 1966).

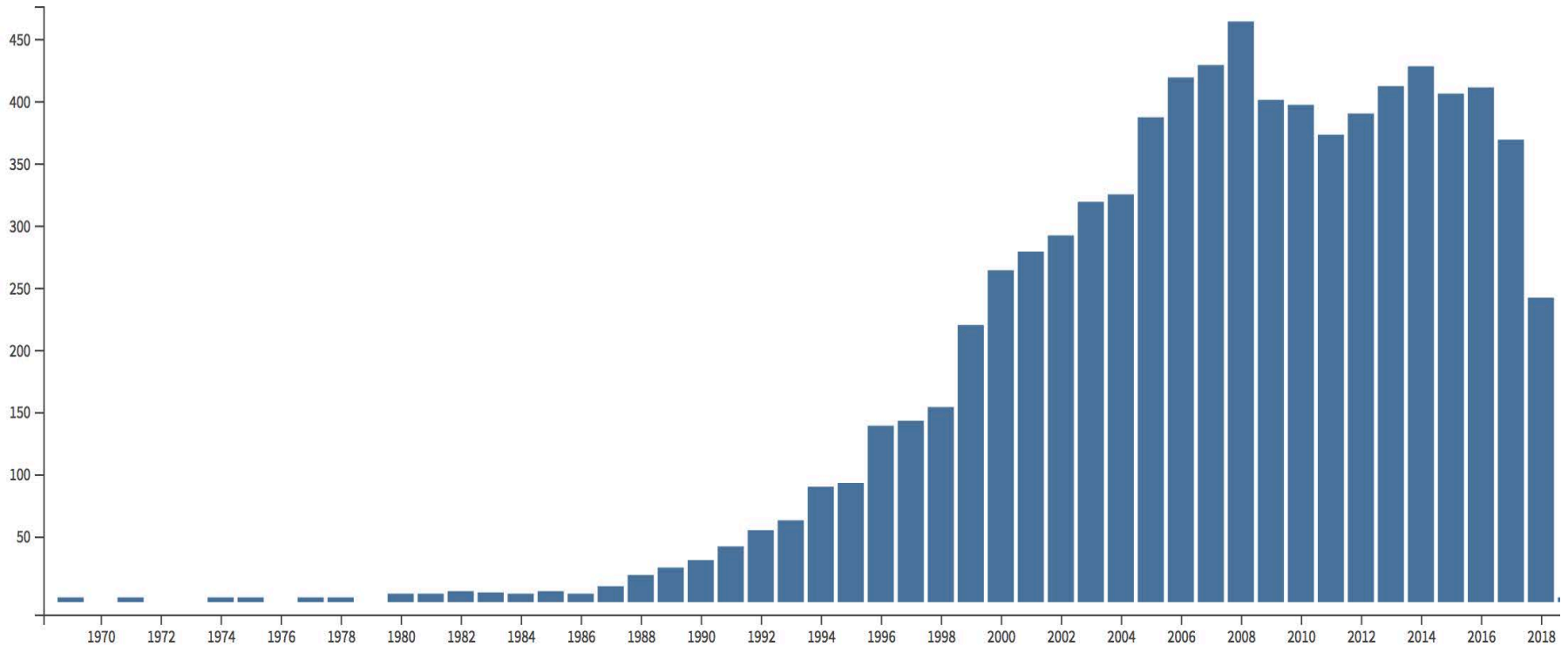
For example, the finite-difference analog of Maxwell's curl equation for the magnetic field:

$$\mathbf{H}^{n+1/2}(\mathbf{r}) = \mathbf{H}^{n-1/2}(\mathbf{r}) - c\Delta t \nabla \times \mathbf{E}^n(\mathbf{r}).$$

Finite-difference Time Domain (FDTD) Method

Yee, K. S., 1966: Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media. IEEE Trans. Antennas Propag., AP-14, 302-307.

As of 10/10/2018: 8,131 citations



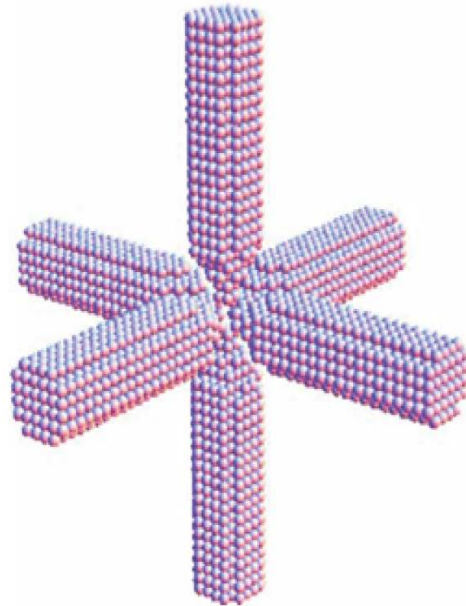
Discrete Dipole Approximation (DDA) Method

(Purcell and Pennypacker 1996; Draine and Flau1994; Yurkin and Hoekstra 2011;...)

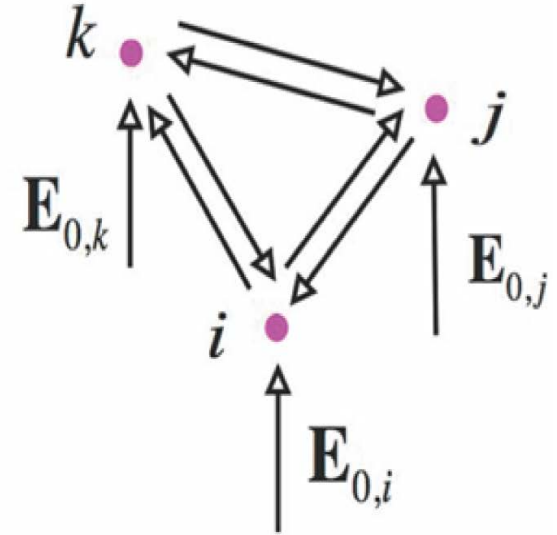


Edward Mills Purcell
Nobel Laureate 1952

(a)



(b)



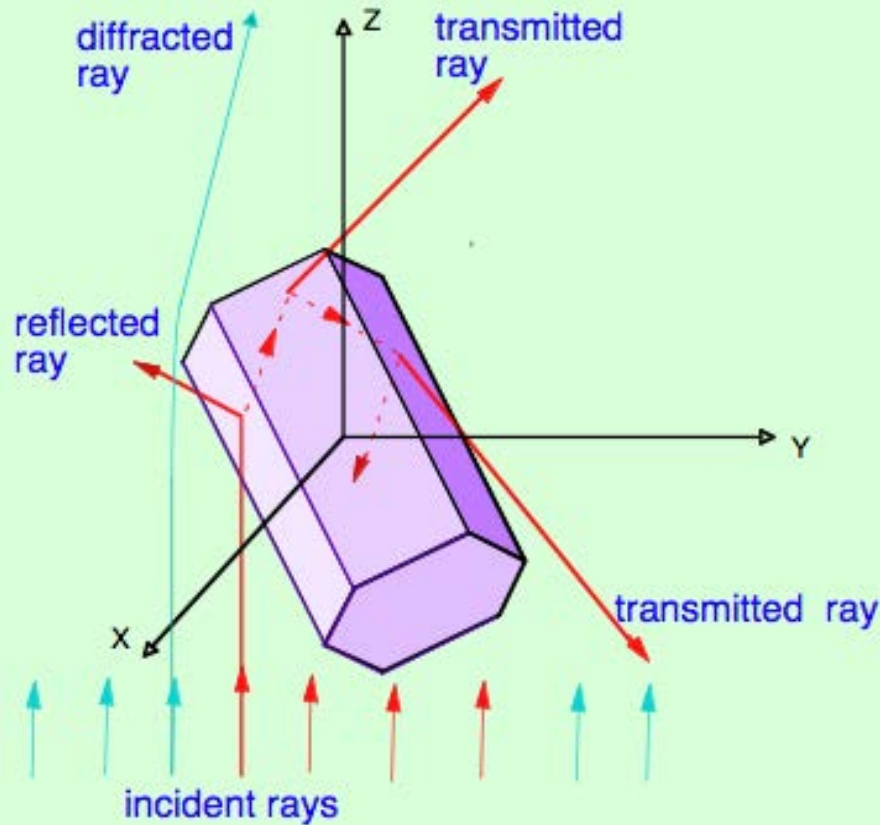
$$\mathbf{P}(\mathbf{r}) = \alpha \mathbf{E}(\mathbf{r})$$

The Clausius-Mossotti (or Lorentz-Lorenz) relation (Lorentz 1880, Lorenz 1880):

$$\alpha = d^3 \frac{3}{4\pi} \frac{m^2 - 1}{m^2 + 2}, d = \text{dipole length}$$

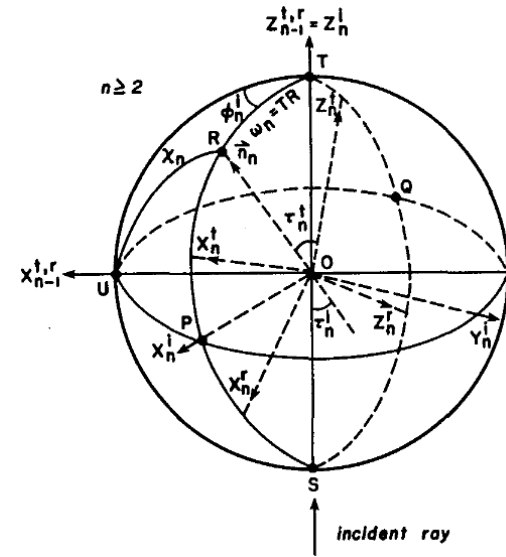
$$\mathbf{P}_i = \alpha_i \left(\mathbf{E}_{0,i} + \sum_{i \neq j} \mathbf{A}_{ij} \cdot \mathbf{P}_j \right) \quad \sigma_{ext} = \frac{4\pi k}{|E_0|^2} \sum_{j=1}^N \text{Im}(\mathbf{E}_{0,j}^* \cdot \mathbf{P}_j)$$

Ray-Tracing Method for Light Scattering Calculation



Conventional Geometric Optics Method

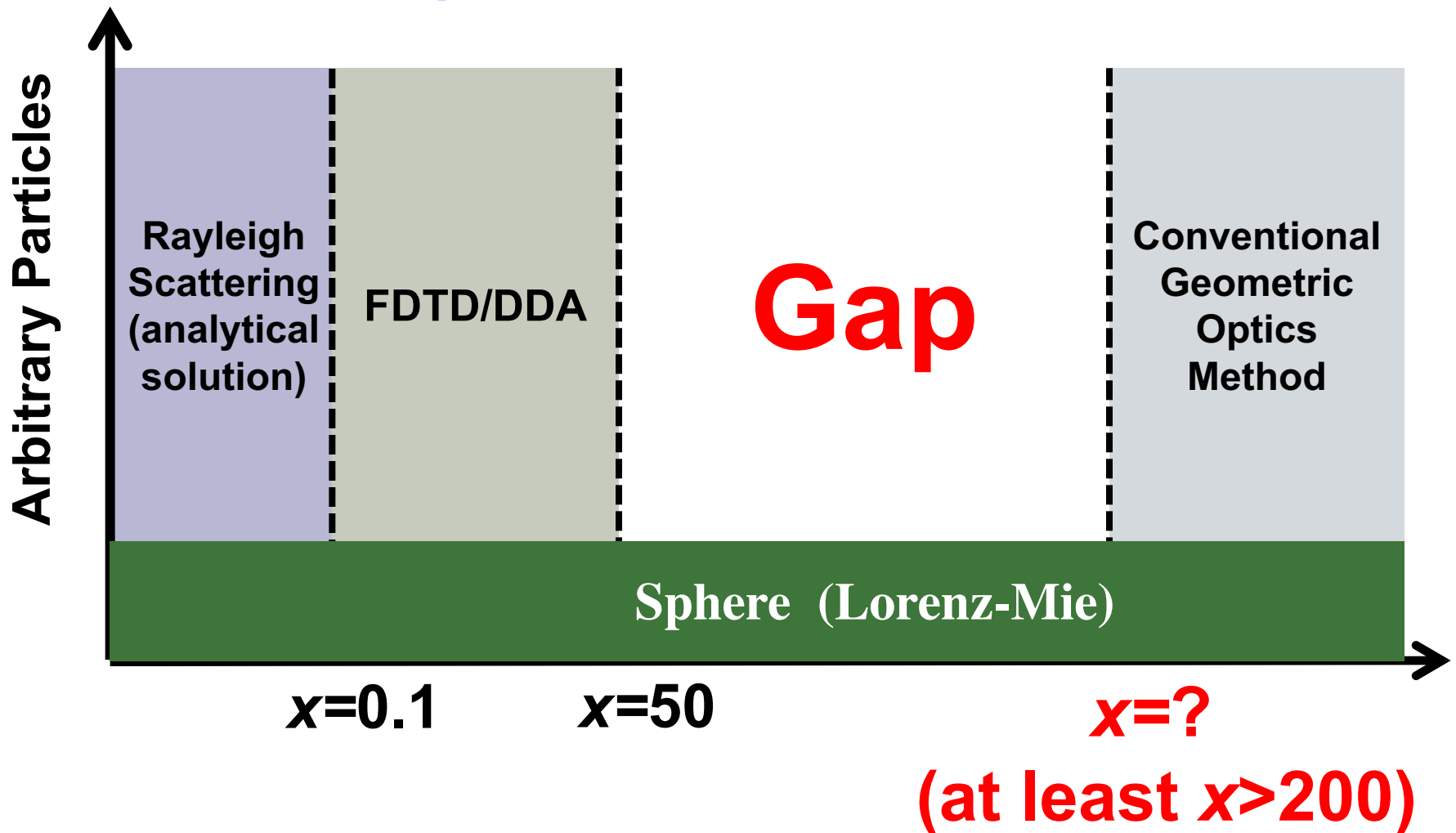
Cai, Q., and K. N. Liou, 1982: Polarized light scattering by hexagonal ice crystals: theory. *Appl. Opt.*, 21, 3569–3580.



- Constant extinction efficiency, 2
- Singularity
- Artificial separation of contributions by diffraction and geometric rays

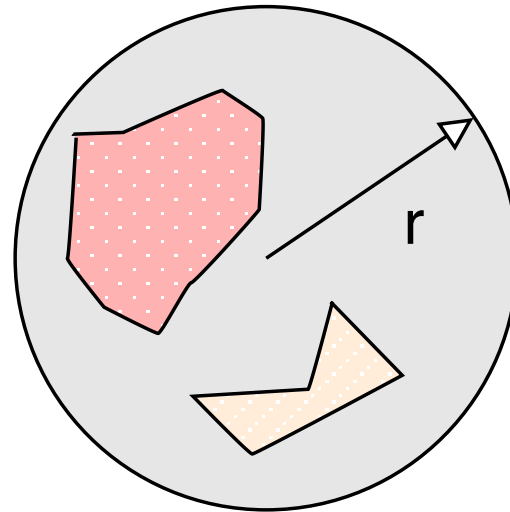
Wendling et al. 1979; Cai and Liou, 1982; Takano and Liou, 1989; Mack 1993; Macke et al. 1996; and many others

Applicability of Light-Scattering Computational Methods

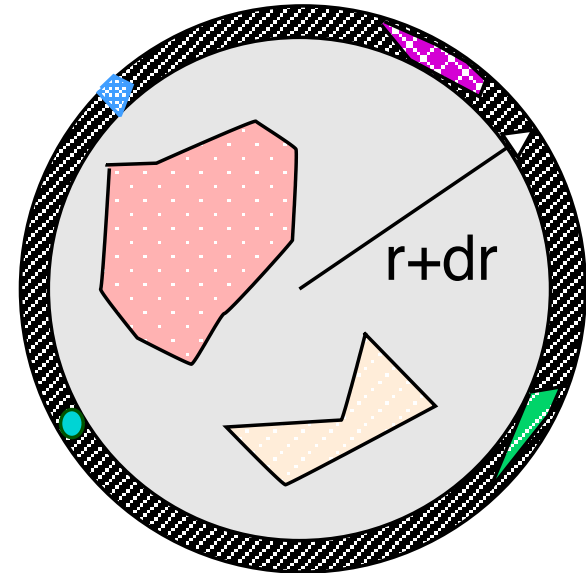


*Breakthrough in Light-
scattering computation*

Invariant Imbedding T-matrix Method (II-TM)



(a)



(b)

*Maxwell's
equations*

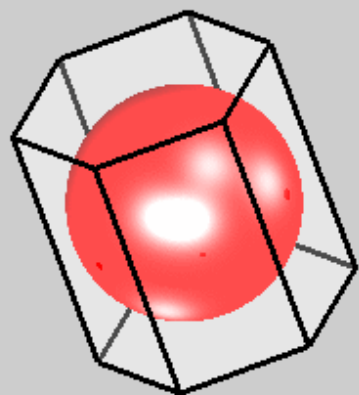


$$\vec{E}(\vec{r}) = \vec{E}_{inc}(\vec{r}) + k^2 \int (m^2 - 1) \vec{G}(\vec{r} - \vec{r}') \cdot \vec{E}(\vec{r}') d^3 \vec{r}'$$

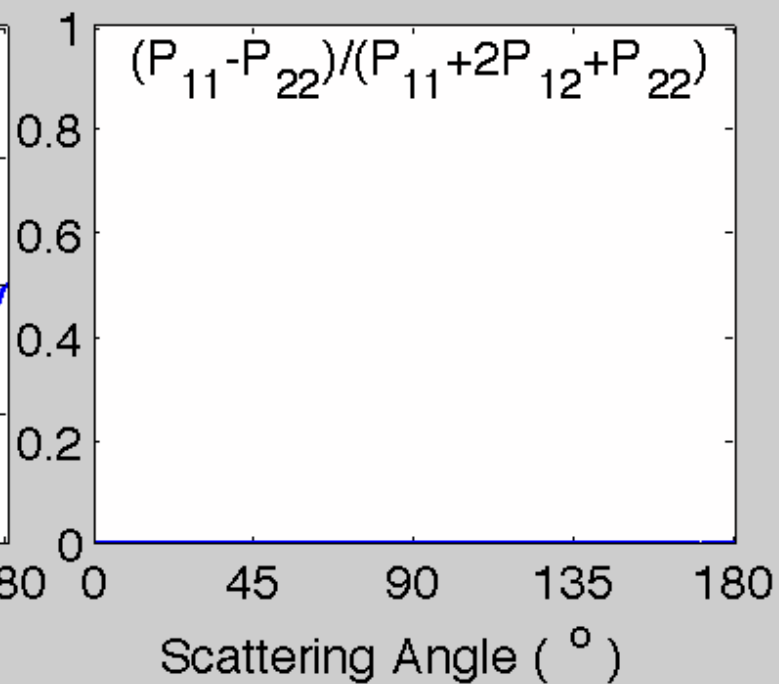
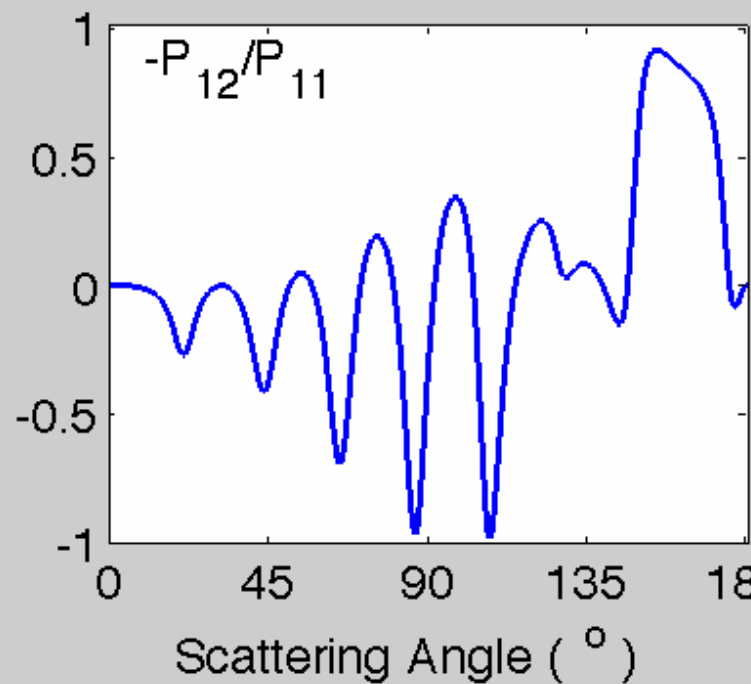
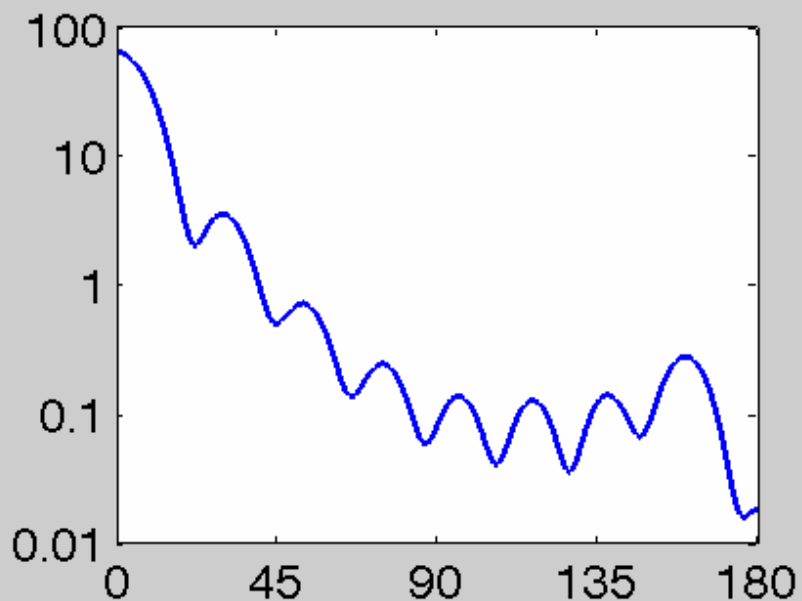
Volume Integral Equation



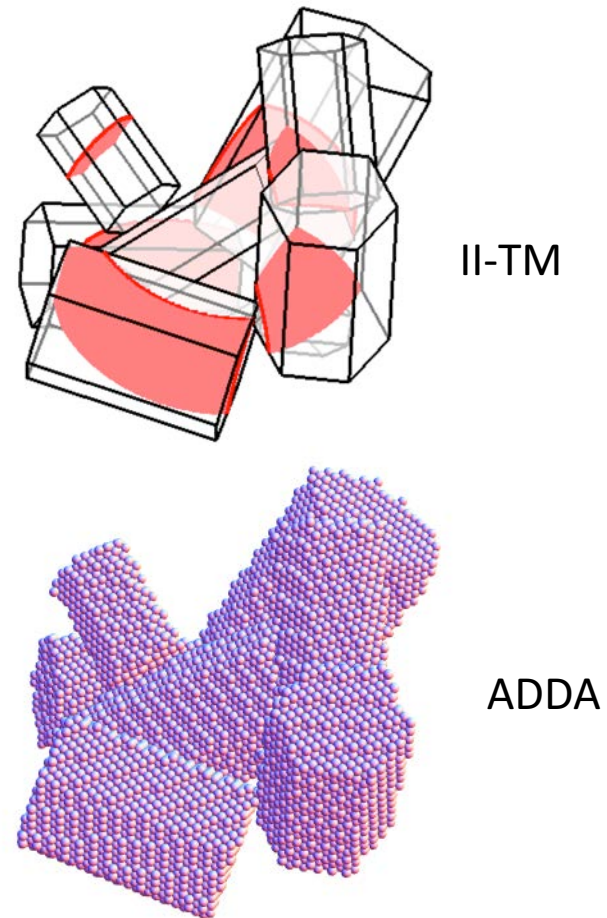
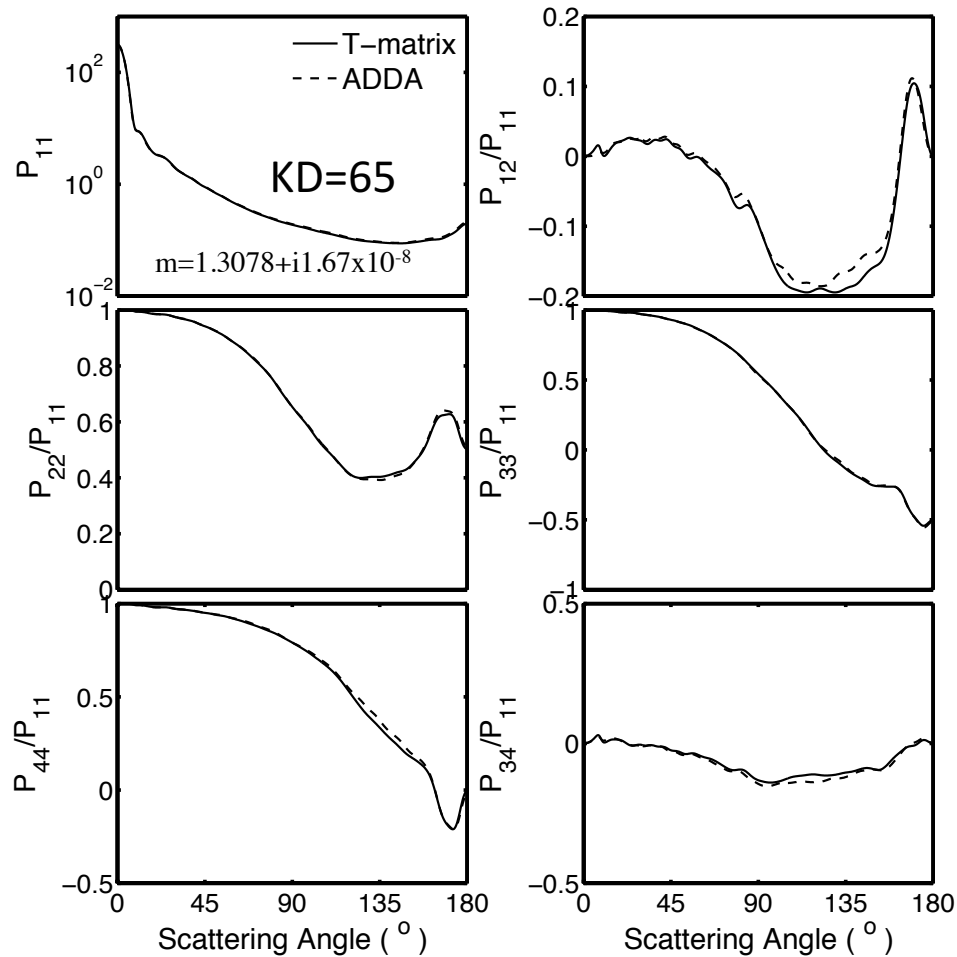
$$T_{mnmn'}(r+dr) = Q_{11}^m(r+dr) + [\mathbf{I} + Q_{12}^m(r+dr)] [\mathbf{I} - T_{mnmn'}(r) Q_{22}^m(r+dr)]^{-1} T_{mnmn'}(r) [\mathbf{I} + \tilde{Q}_{12}^m(r+dr)]$$



P_{11}

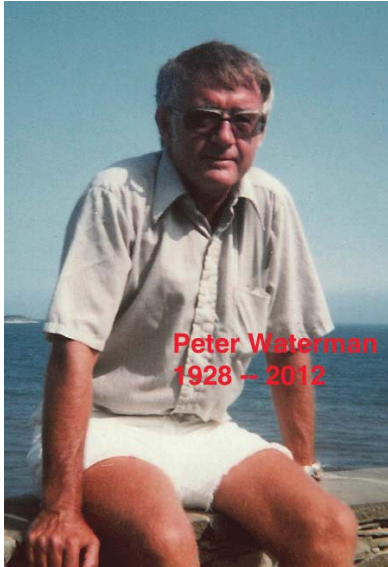


Validation



In the ADDA simulation, 1056 orientations with 128 scattering planes are set to achieve the randomness. Bi and Yang (2014).

II-TM is different from the conventional T-matrix method (Extended Boundary Condition Method, EBCM)



Peter Waterman
1928 – 2012

EBCM:

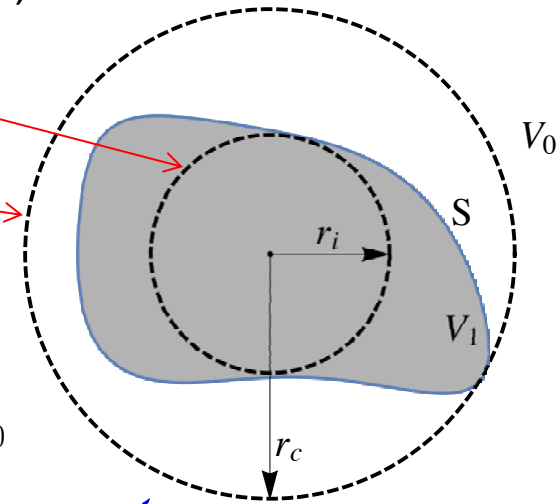
Waterman PC, 1965: Matrix formulation of electromagnetic scattering. Proc. IEEE 53, 805-12

Waterman PC, 1971: Symmetry, unitary, and geometry in electromagnetic scattering. Phys Rev D 3, 825-39

(Mishchenko and Martin, JQSRT, 123:2-7, 2013)

Inscribed sphere

Circumscribing sphere



$$\vec{E}^{\text{inc}}(\vec{r}') = - \int_S ds \{ i\omega\mu_0 [\hat{n} \times \vec{H}(\vec{r})] \cdot \vec{G}(\vec{r}, \vec{r}') + [\hat{n} \times \vec{E}(\vec{r})] \cdot [\nabla \times \vec{G}(\vec{r}, \vec{r}')] \}, \quad \vec{r}' \in V_1$$

$$\vec{E}^{\text{sca}}(\vec{r}') = \int_S ds \{ i\omega\mu_0 [\hat{n} \times \vec{H}(\vec{r})] \cdot \vec{G}(\vec{r}, \vec{r}') + [\hat{n} \times \vec{E}(\vec{r})] \cdot [\nabla \times \vec{G}(\vec{r}, \vec{r}')] \}, \quad \vec{r}' \in V_0$$

$$T = -RgQ[Q]^{-1}$$

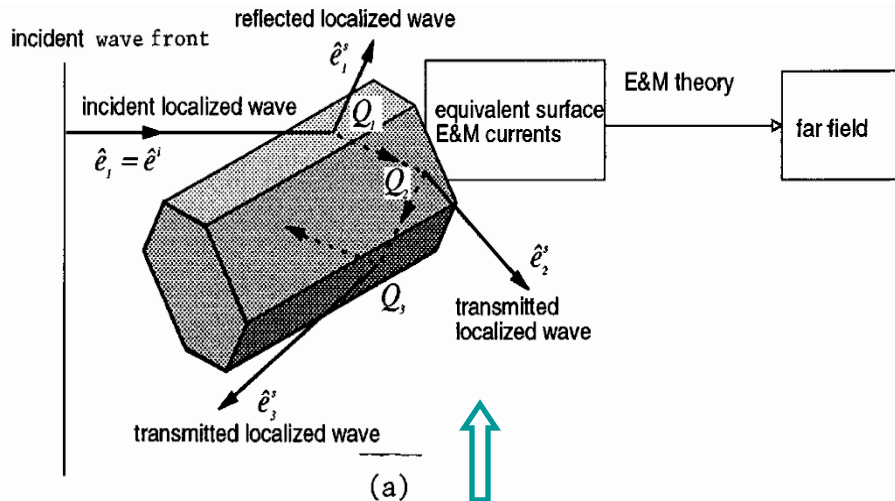
Surface Integral Equations

For technical details, please see Mishchenko MI, Travis LD and Lacis AA, *Scattering, Absorption, and Emission of Light by Small Particles*. Cambridge: Cambridge University Press; 2002.

Physical-Geometric Optics Method (PGOM)

Yang and Liou (1996)

PGOMS – Surface-integral equation based



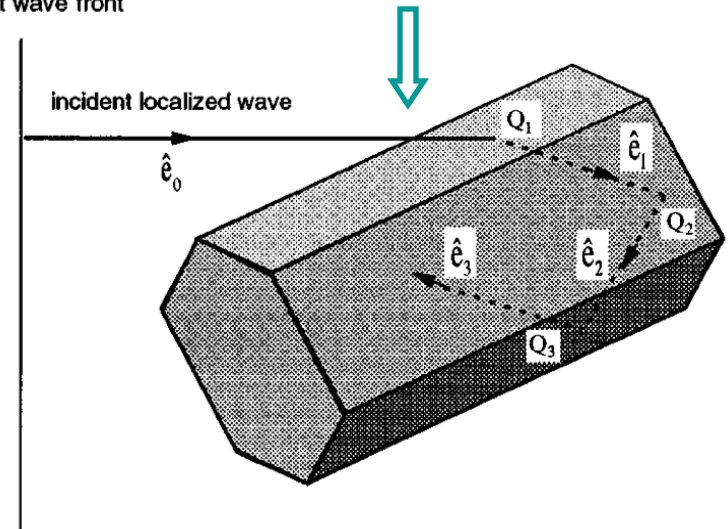
$$\mathbf{E}_s(\mathbf{r})|_{kr \rightarrow \infty} = \frac{\exp(ikr)}{-ikr} \frac{k^2}{4\pi} \mathbf{n} \times \iint_S \{ \mathbf{n}_S \times \mathbf{E}(\mathbf{r}') - \mathbf{n} \times [\mathbf{n}_S \times \mathbf{H}(\mathbf{r}')] \} \times \exp(-ik\mathbf{n} \cdot \mathbf{r}') d^2\mathbf{r}'$$

Yang and Liou (1997)

PGOMV – Volume-integral equation based

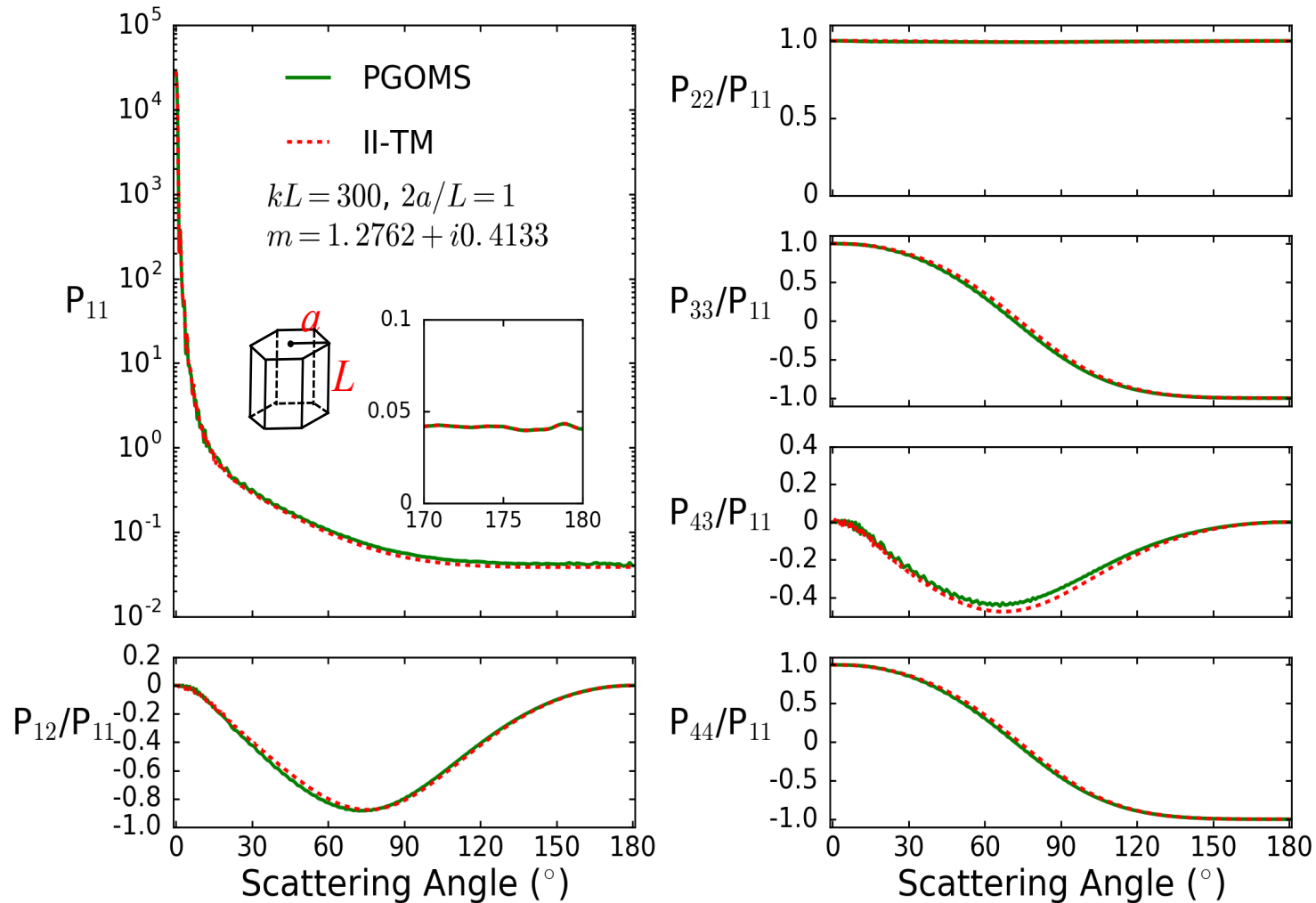
$$\mathbf{E}_s(\mathbf{r}) = \frac{k^2 \exp(ikr)}{4\pi r} \iint_V [\varepsilon(\mathbf{r}') - 1] \{ \mathbf{E}(\mathbf{r}') - \mathbf{n}[\mathbf{n} \cdot \mathbf{E}(\mathbf{r}')] \} \exp(-ik\mathbf{n} \cdot \mathbf{r}') d^3\mathbf{r}'$$

incident wave front



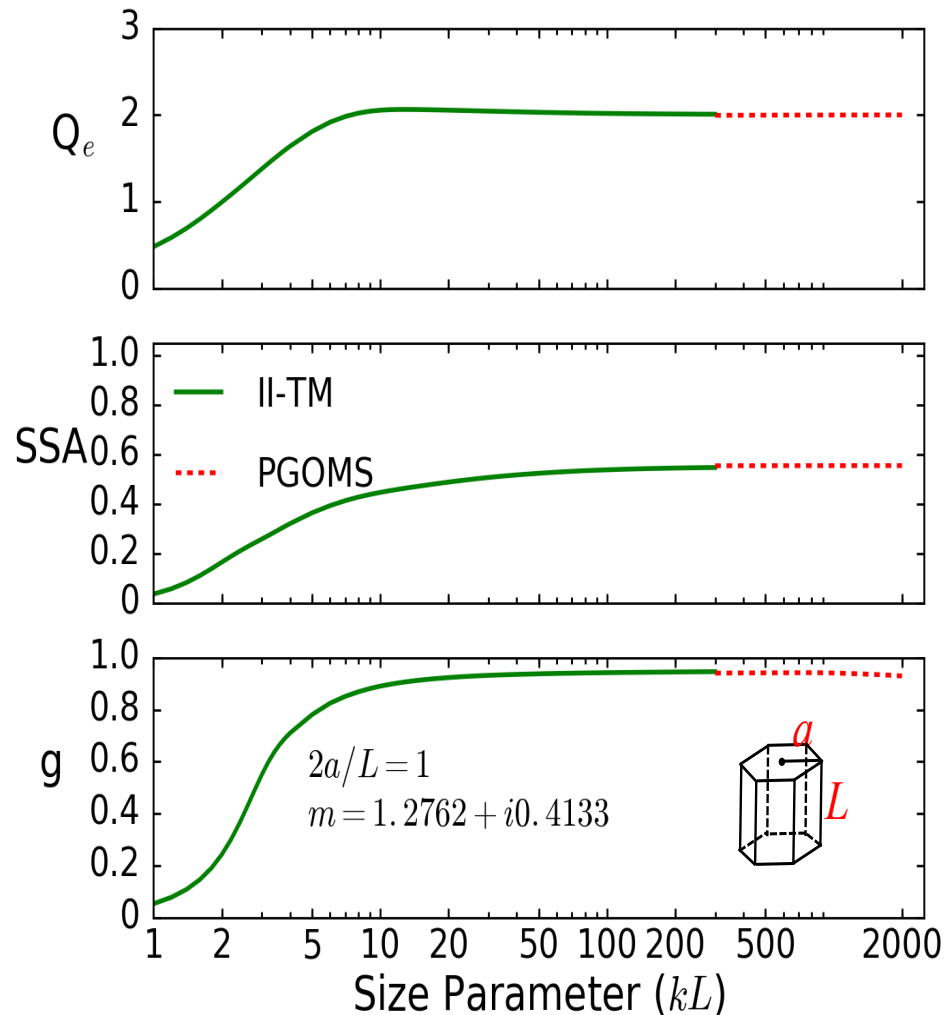
New improvements by our research group at Texas A&M University using computer graphics techniques

PGOMS vs II-TM



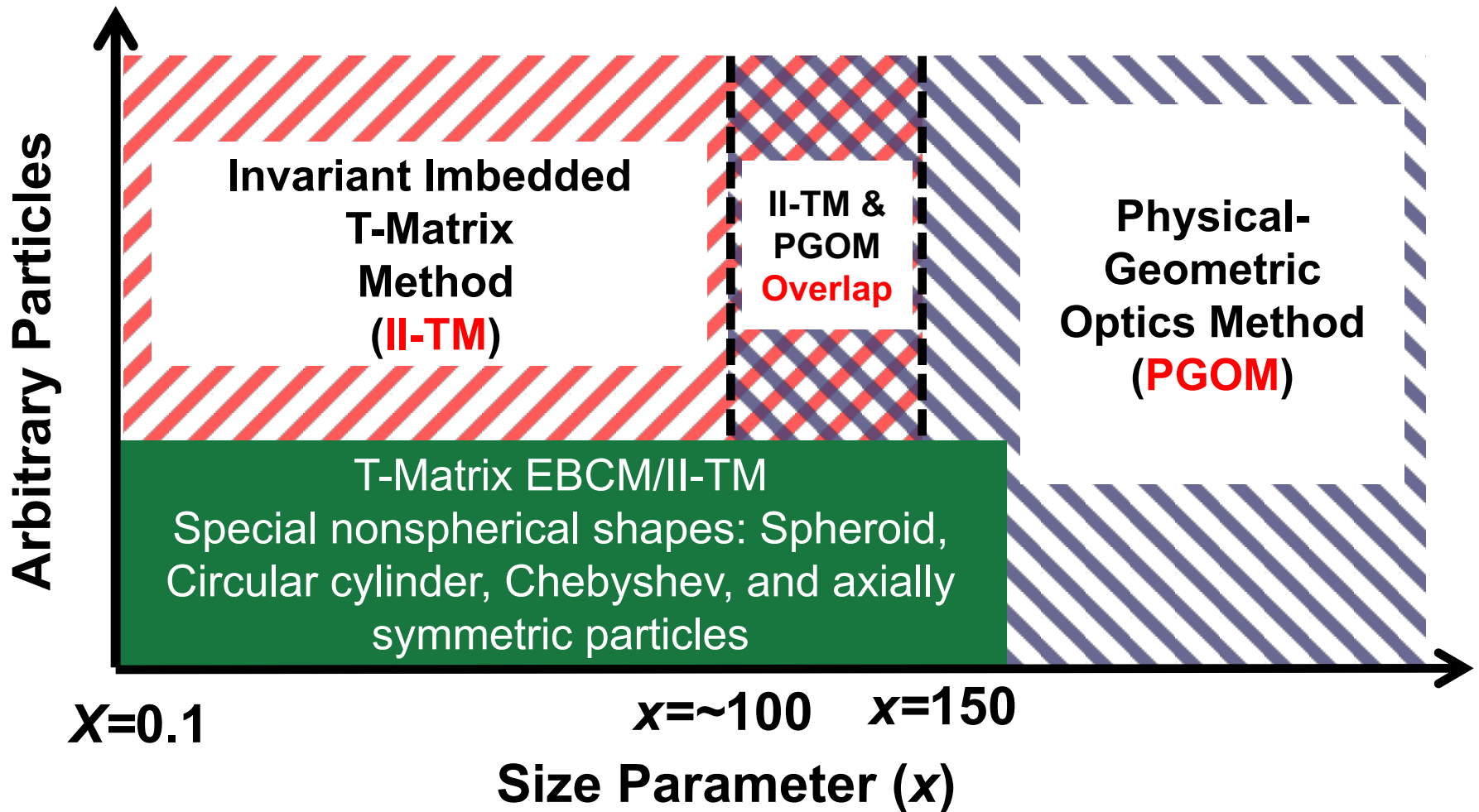
Comparison of the phase matrix elements computed by PGOMS and IITM. The particle is a hexagonal column with aspect ratio 1. The refractive index is $1.2762+i0.4133$, the ice refractive index at $12\mu\text{m}$ wavelength. The inset plots show the P_{11} element for 170° - 180° scattering angles. The size parameter is $kL=300$, or $ka=150$.

PGOMS vs II-TM



Extinction efficiency (Q_e), single-scattering (SSA), and asymmetry factor (g) computed by II-TM and PGOM. The particle is a hexagonal column with aspect ratio 1. The refractive index is $1.2762 + i0.4133$ that is the ice refractive index at $12 \mu m$ wavelength

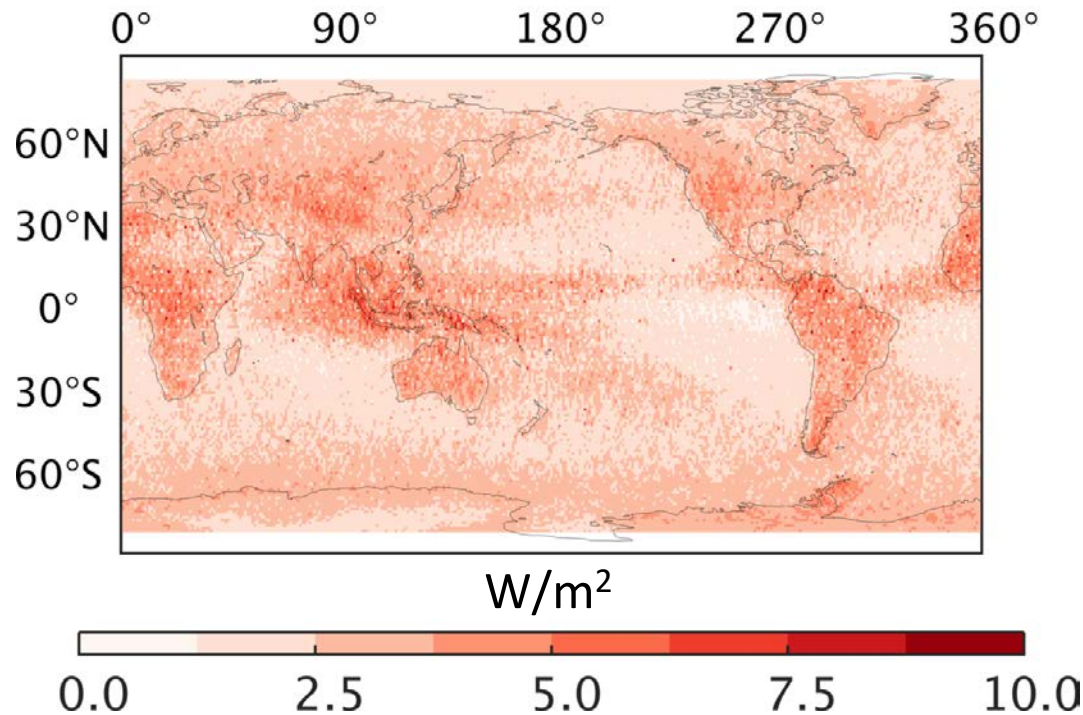
Breakthrough: A combination of II-TM and PGOM can accurately cover the entire size parameter region



*Scattering effect is important
in radiative transfer
simulation even in the
infrared region*

*The current **RRTM-G** neglects
the scattering effect in LW
bands*

TOA Upward Flux Biases



Areas containing large biases

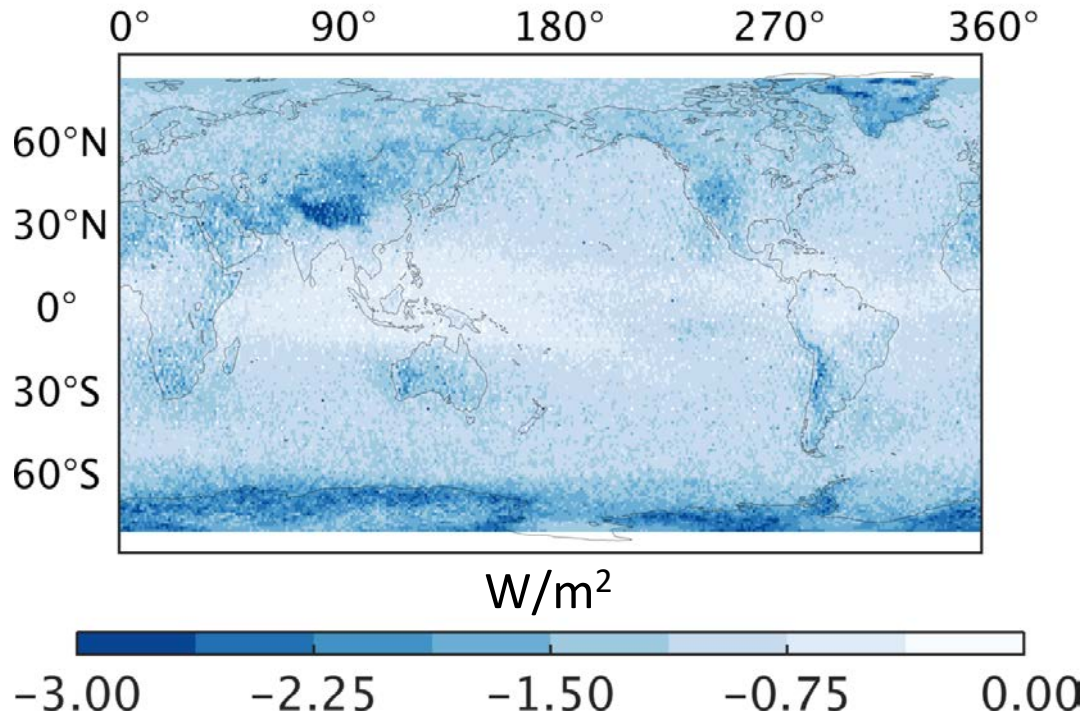
- **Intertropical Convergence Zone (ITCZ)**
- **Pacific warm pool**
- **Tibetan Plateau**

Large biases (up to 12 W/m²)

Positive biases mean that the TOA upward fluxes are overestimated when LW scattering is ignored

Kuo, C.-P., P. Yang, X. Huang, D. Feldman, M. Flanner, C. Kuo, and E. J. Mlawer, 2017: *Journal of Advances in Modeling Earth Systems*.

Surface Downward Flux Biases



Areas containing large biases

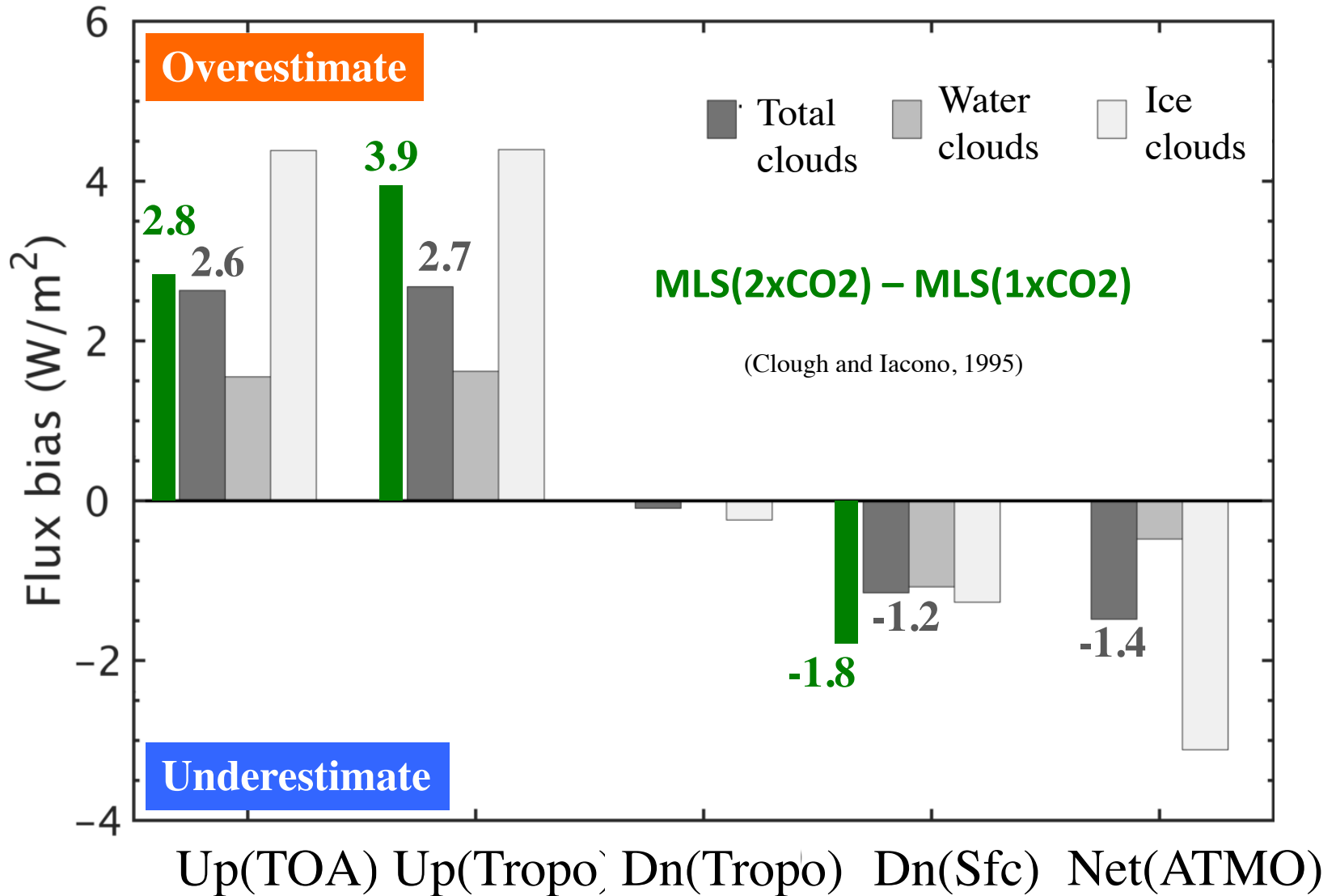
- Dry and high regions
- Tibetan Plateau
- Antarctic
- Greenland

Large biases ($\sim -3.6 \text{ W/m}^2$)

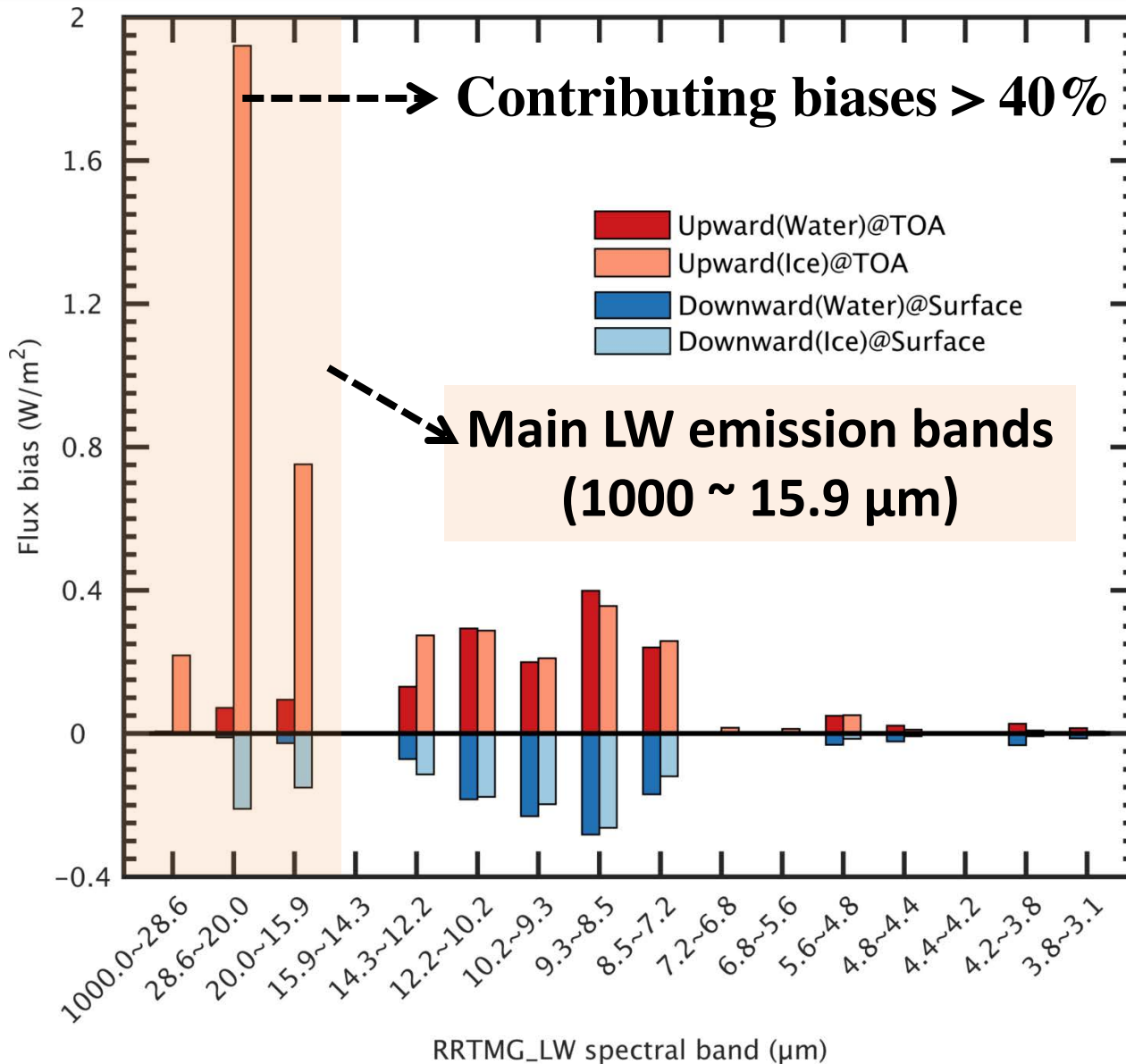
Negative biases mean that the surface downward fluxes are underestimated when LW scattering is ignored

Kuo, C.-P., P. Yang, X. Huang, D. Feldman, M. Flanner, C. Kuo, and E. J. Mlawer, 2017: *Journal of Advances in Modeling Earth Systems*.

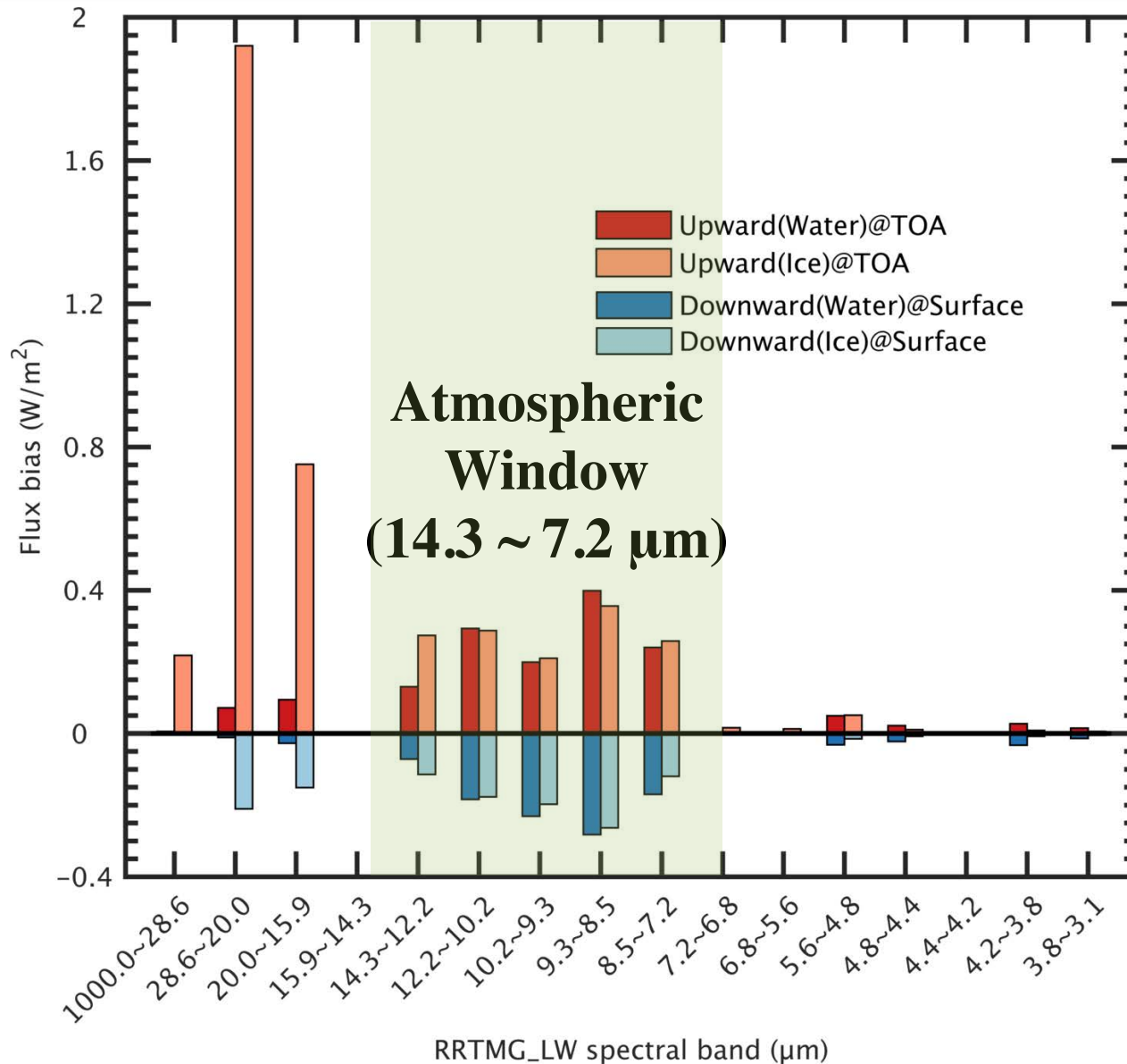
Flux Biases



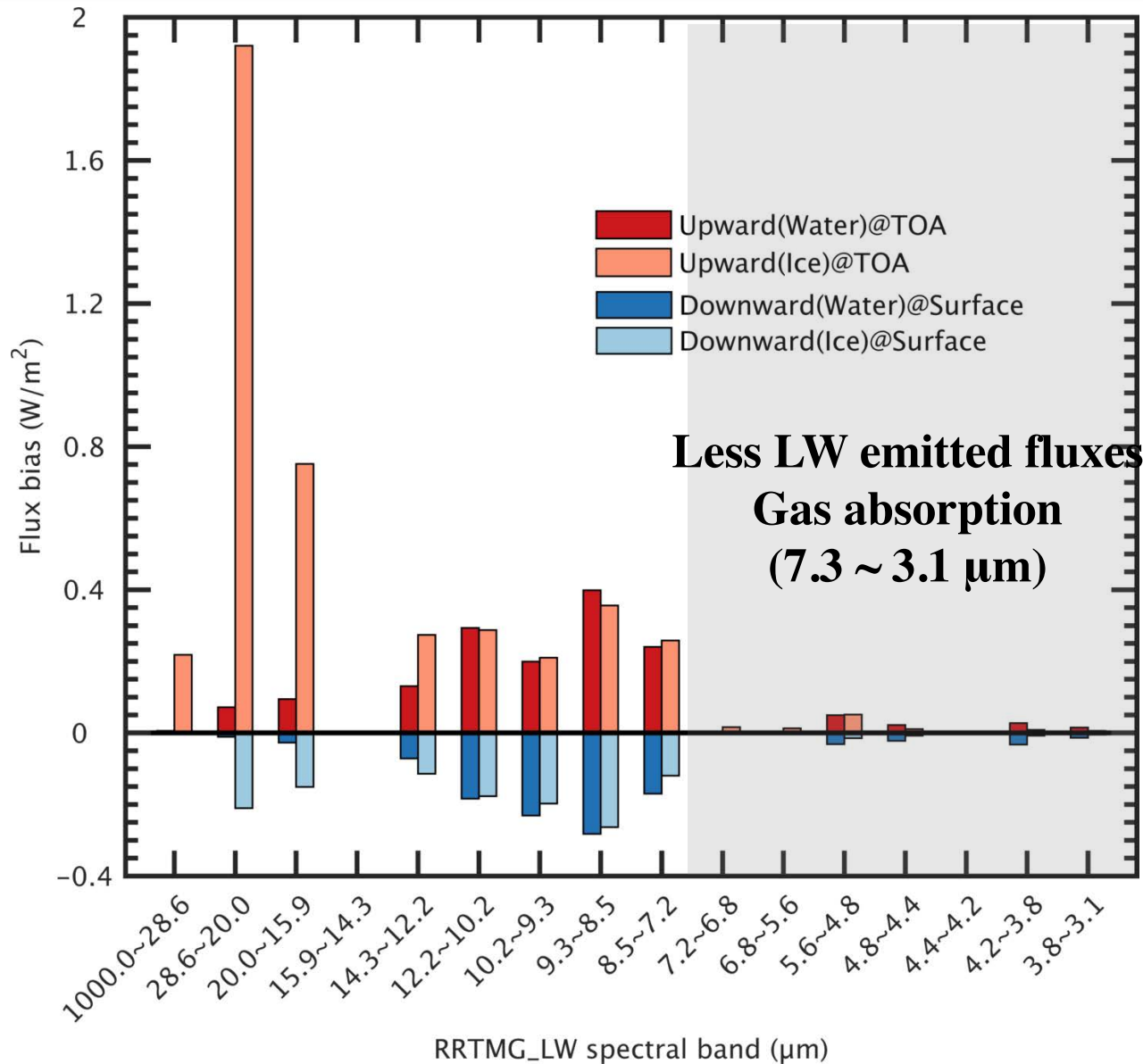
Spectral Analyses



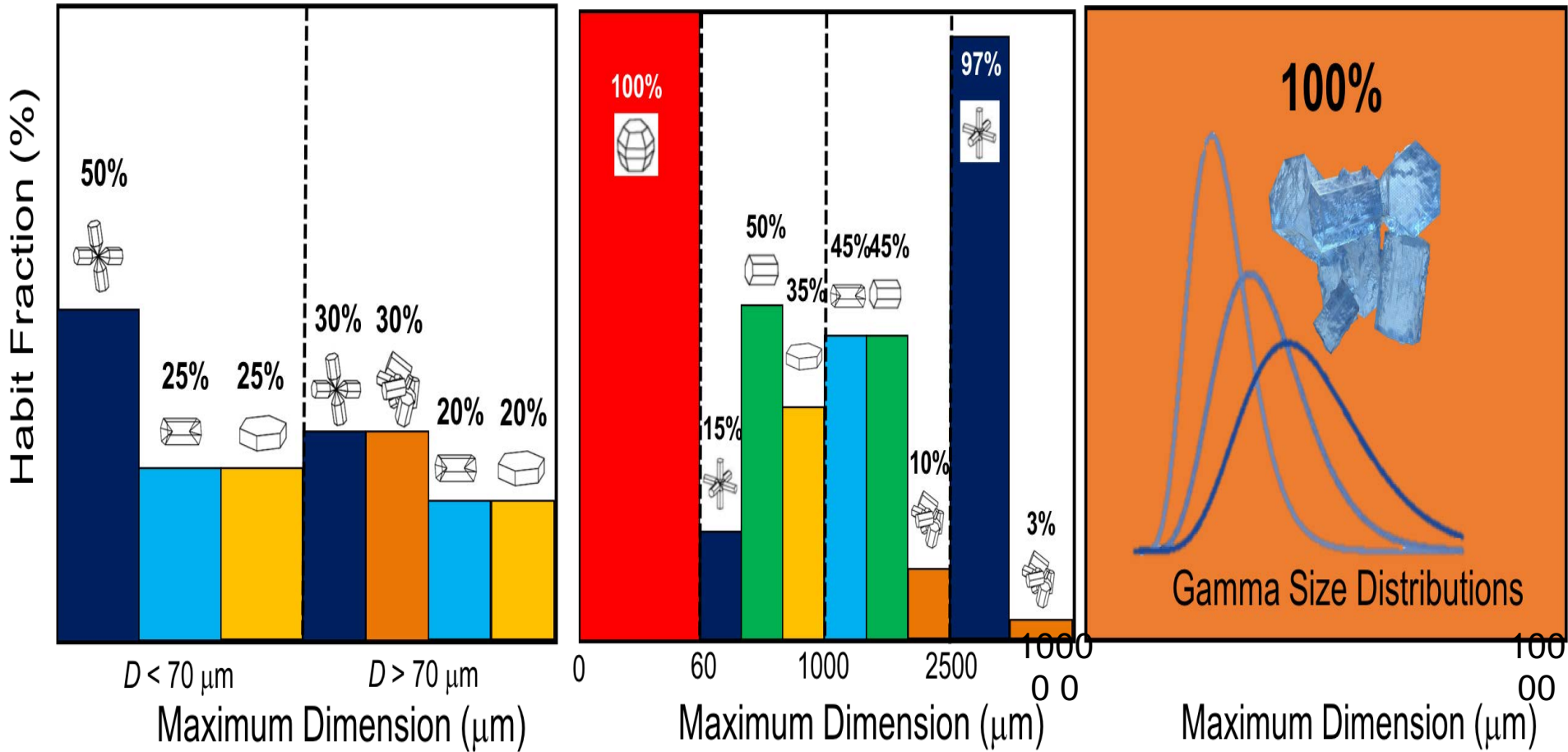
Spectral Analyses



Spectral Analyses

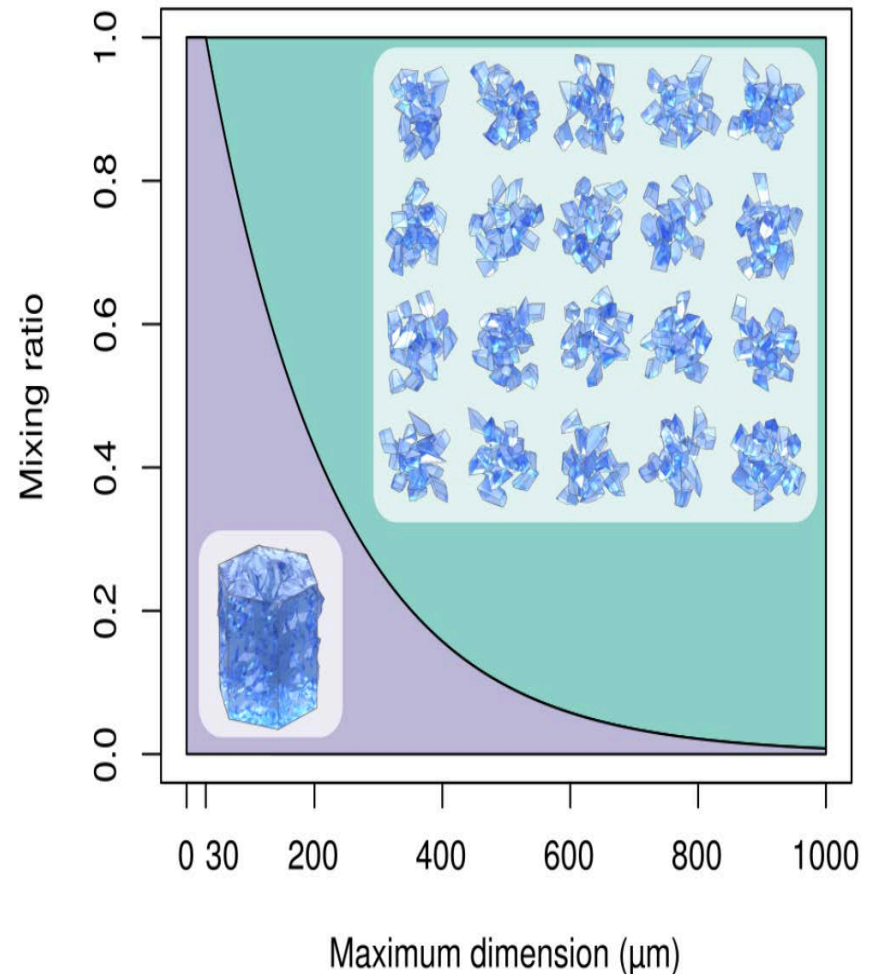
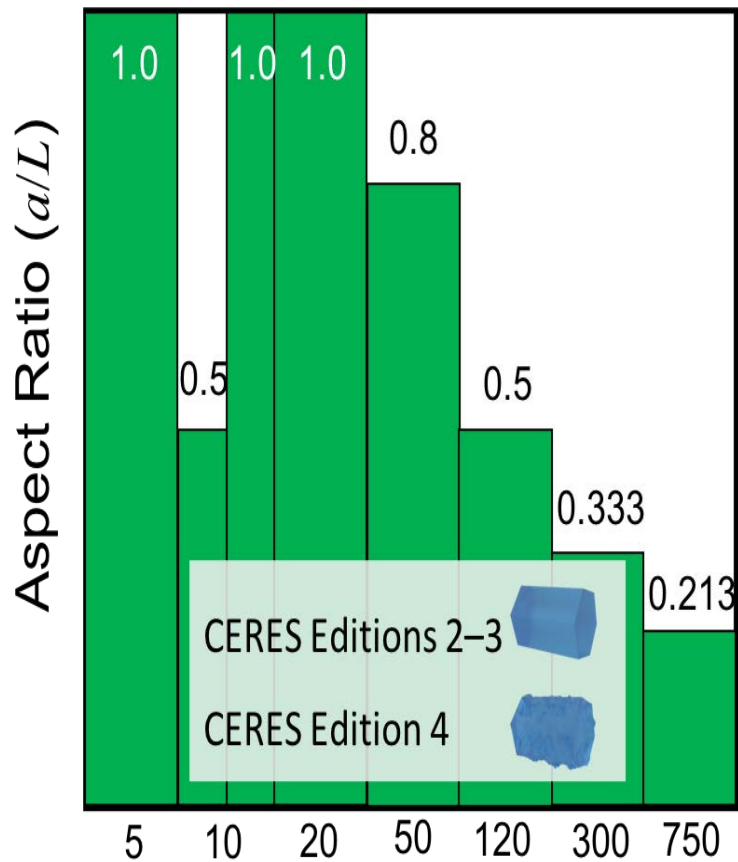


MODIS Ice Particle Models (Collections 4, 5, 6)



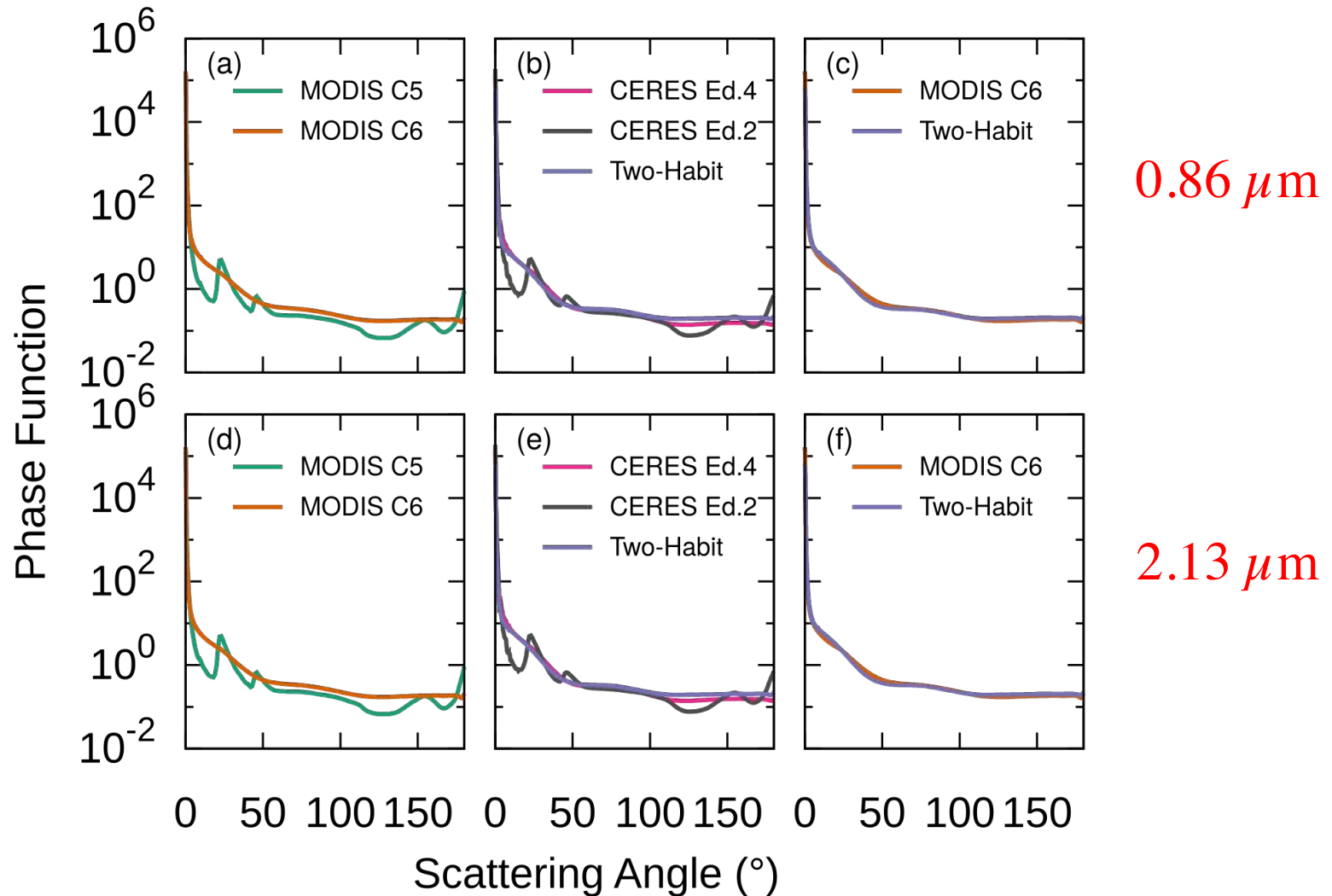
References: King et al. 2004, Baum et al. 2005, Platnick et al. 2017

CERES ice particle models (Editions 2-4, and a two-habit model for future Edition 5)



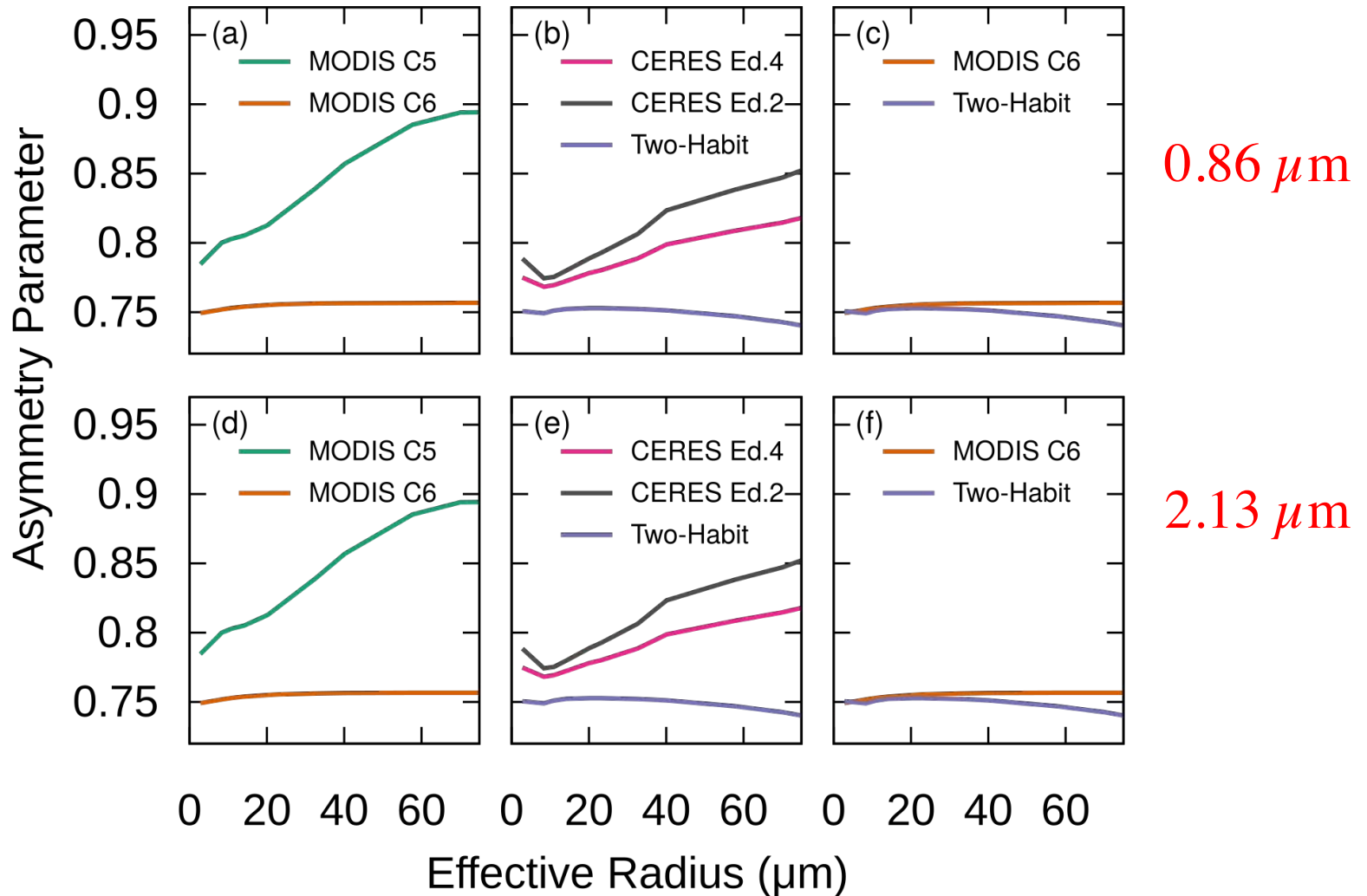
References: Minnis et al. 1993, 2011; Loeb et al. 2017

Phase Function Comparison



Comparison of the phase functions at wavelength $0.86 \mu\text{m}$ based on (a) MODIS Collection 4, 5, and 6; (b) CERES Edition 2, 4, and the Two-habit model; and (c) MODIS Collection 6 and the Two-habit model. Diagrams (d), (e), and (f) in bottom rows are counterpart of (a), (b), (c) at wavelength $2.13 \mu\text{m}$. Effective radius is fixed at $30 \mu\text{m}$.

Asymmetry Factor Comparison



Comparison of the asymmetry parameter at wavelength $0.86 \mu\text{m}$ based on (a) MODIS Collection 4, 5, and 6; (b) CERES Edition 2, 4, and the Two-habit model; and (c) MODIS Collection 6 and the Two-habit model. Diagrams (d), (e), and (f) in bottom rows are counterpart of (a), (b), (c) at wavelength $2.13 \mu\text{m}$. Effective radius is fixed at $30 \mu\text{m}$.

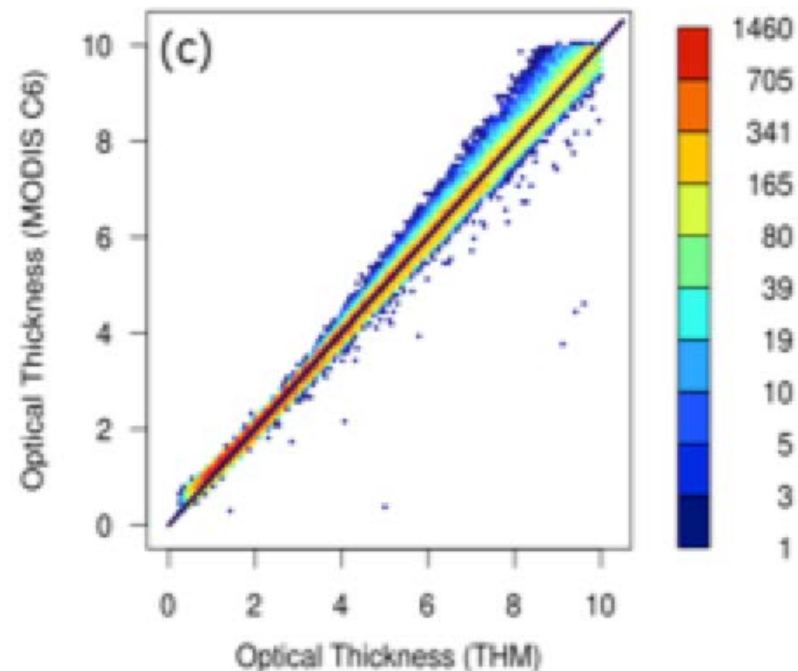
The similarity relation at a non-absorptive wavelength (van de Hulst 1971, 1974)

$$(1-g) \tau = (1-g') \tau'$$

Thus

$g_{C6} \cong g_{THM}$ leads to $\tau_{C6} \cong \tau_{THM}$

$g_{C5} > g_{C6}$ leads to $\tau_{C5} > \tau_{C6}$

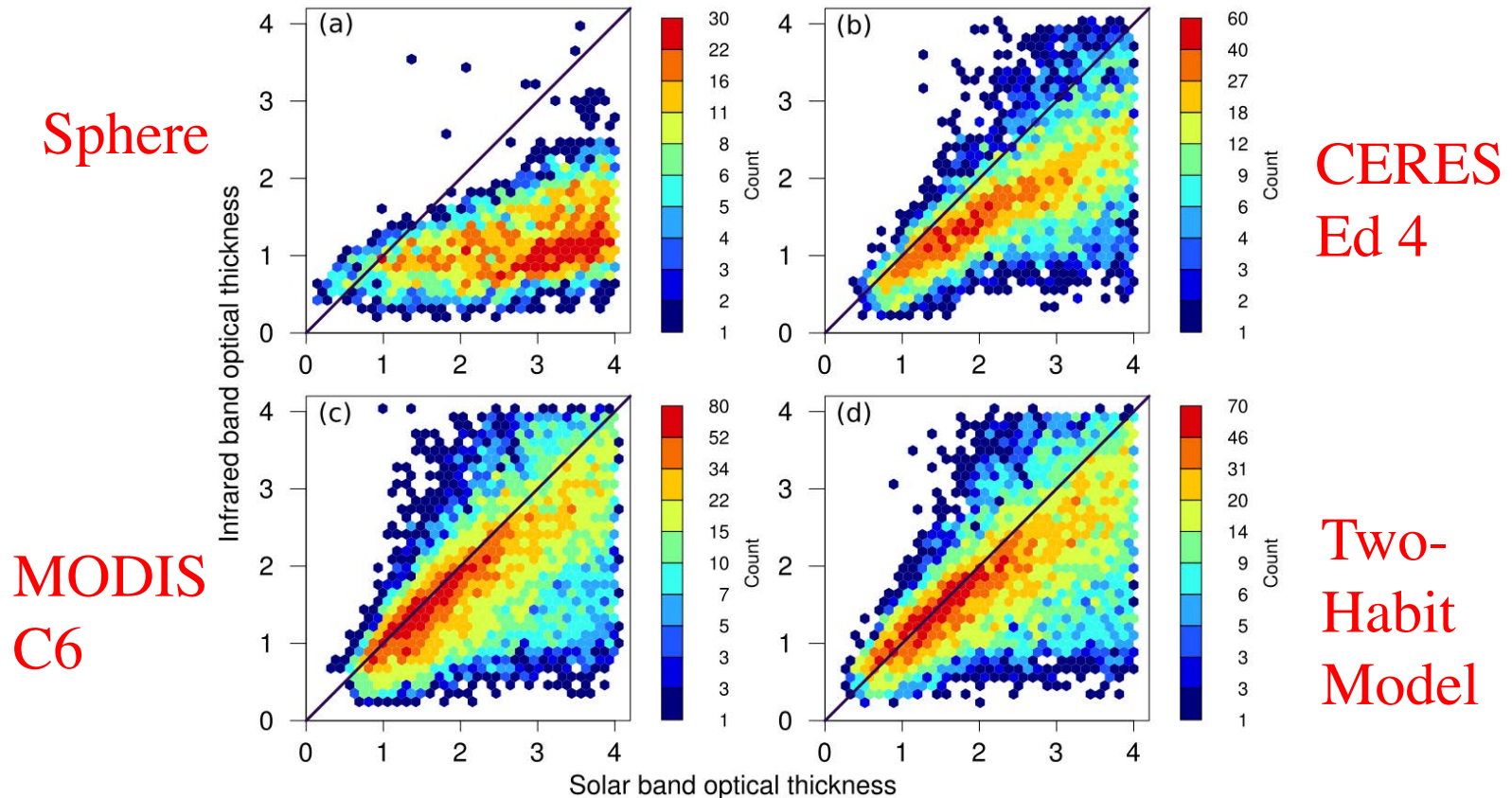


Spectral consistency

- *Nakajima-King bispectral method based on two solar bands*
- *Split window technique based on thermal infrared bands*

Retrievals based on the two methods should be consistent!

Consistency Check (VIS-NIR vs IR retrieval techniques)

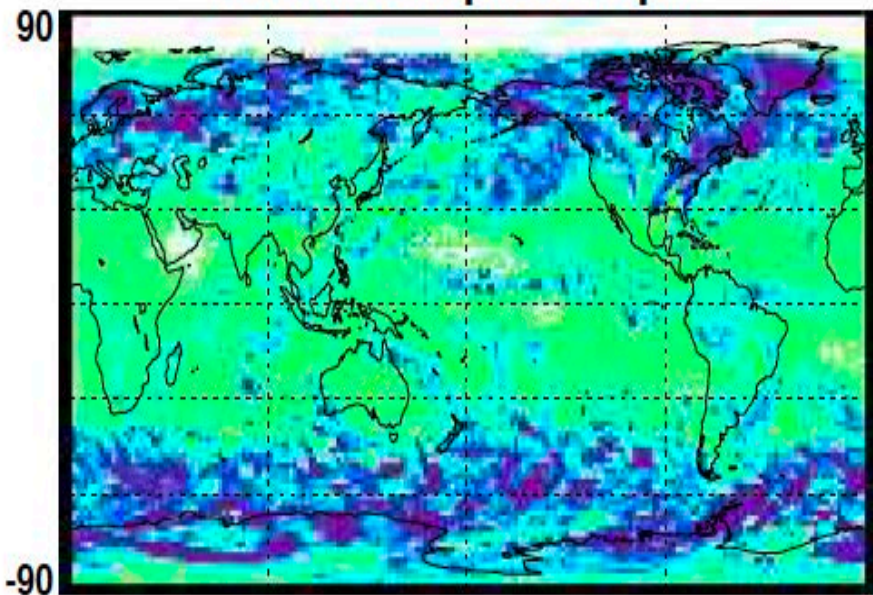


Comparison of retrieved optical thickness values from a shortwave method (the Nakajima-King bi-spectral method) and a longwave method (the split-window technique). (a) Ice ³⁶ sphere, (b) CERES Edition 4 model, (c) MODIS Collection 6

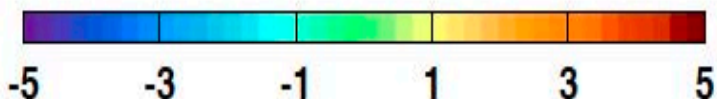
Loeb et al., 2018: Impact of ice microphysics on satellite cloud retrievals and broadband flux radiative transfer model calculations. *J. Climate*, 31, 1851-1864.

Cloud Property Differences at Aqua Overpass Time (THM minus Smooth)

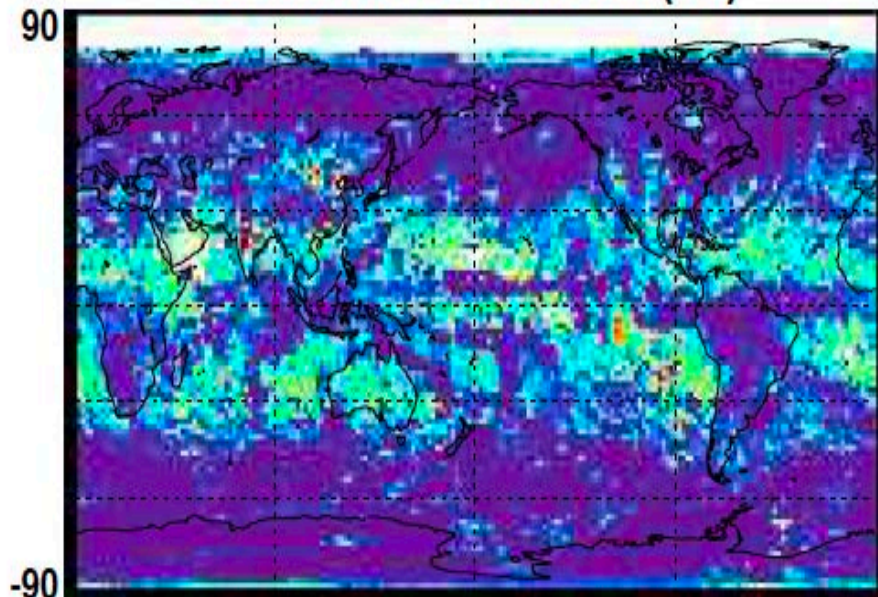
Cloud Ice Optical Depth



Cloud Optical Depth Difference (Wm^{-2})



Cloud Ice Particle Size(Re)



Effective Radius Difference (%)

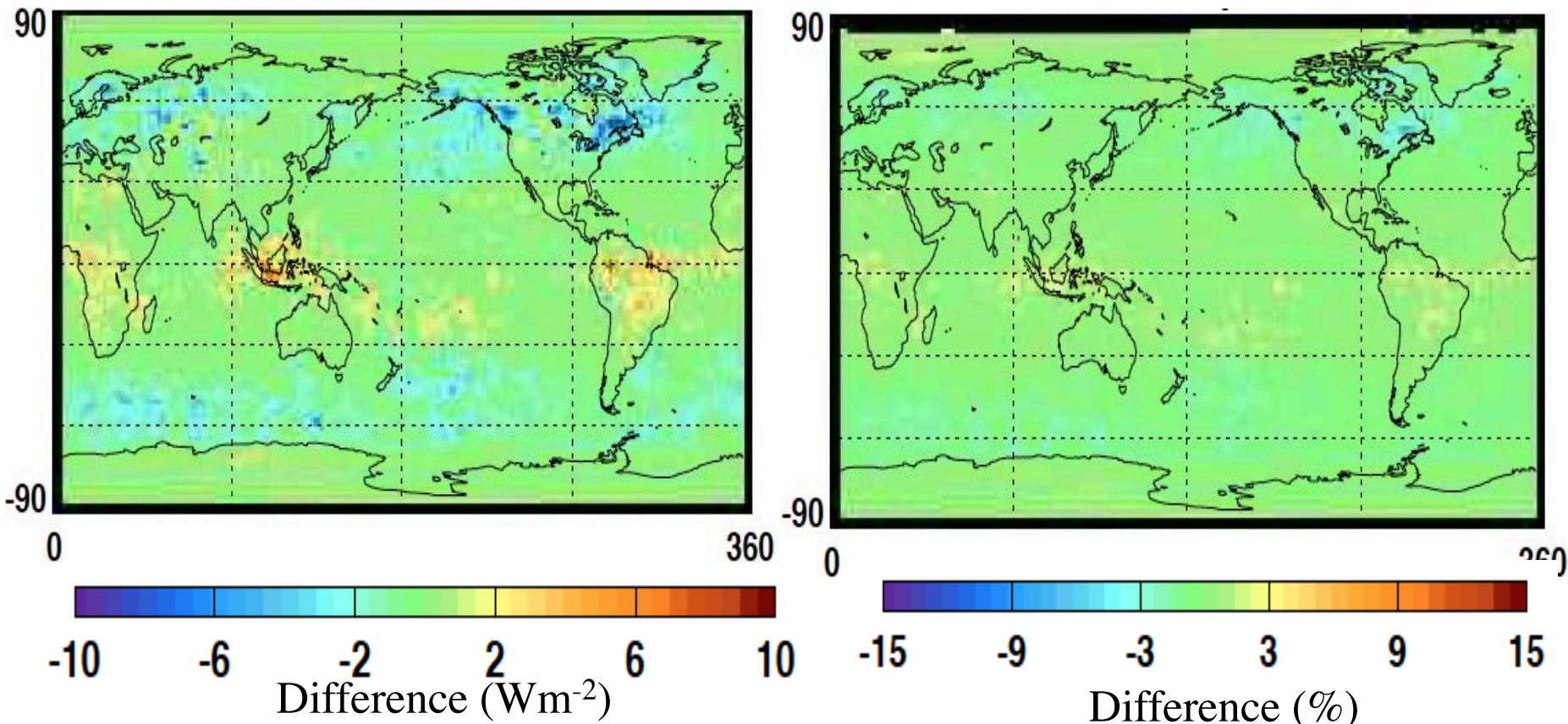


- Overall optical depth difference is -2.3 (-28% of Global Mean) and RMS difference is 2.8 (32% of GM).
- Overall effective radius difference is -3.9 μm (16% of GM) and RMS difference is 5.2 μm (16% of GM).

Loeb et al., 2018: Impact of ice microphysics on satellite cloud retrievals and broadband flux radiative transfer model calculations. *J. Climate*, 31, 1851-1864.

SW TOA Flux Difference at Aqua Overpass Time

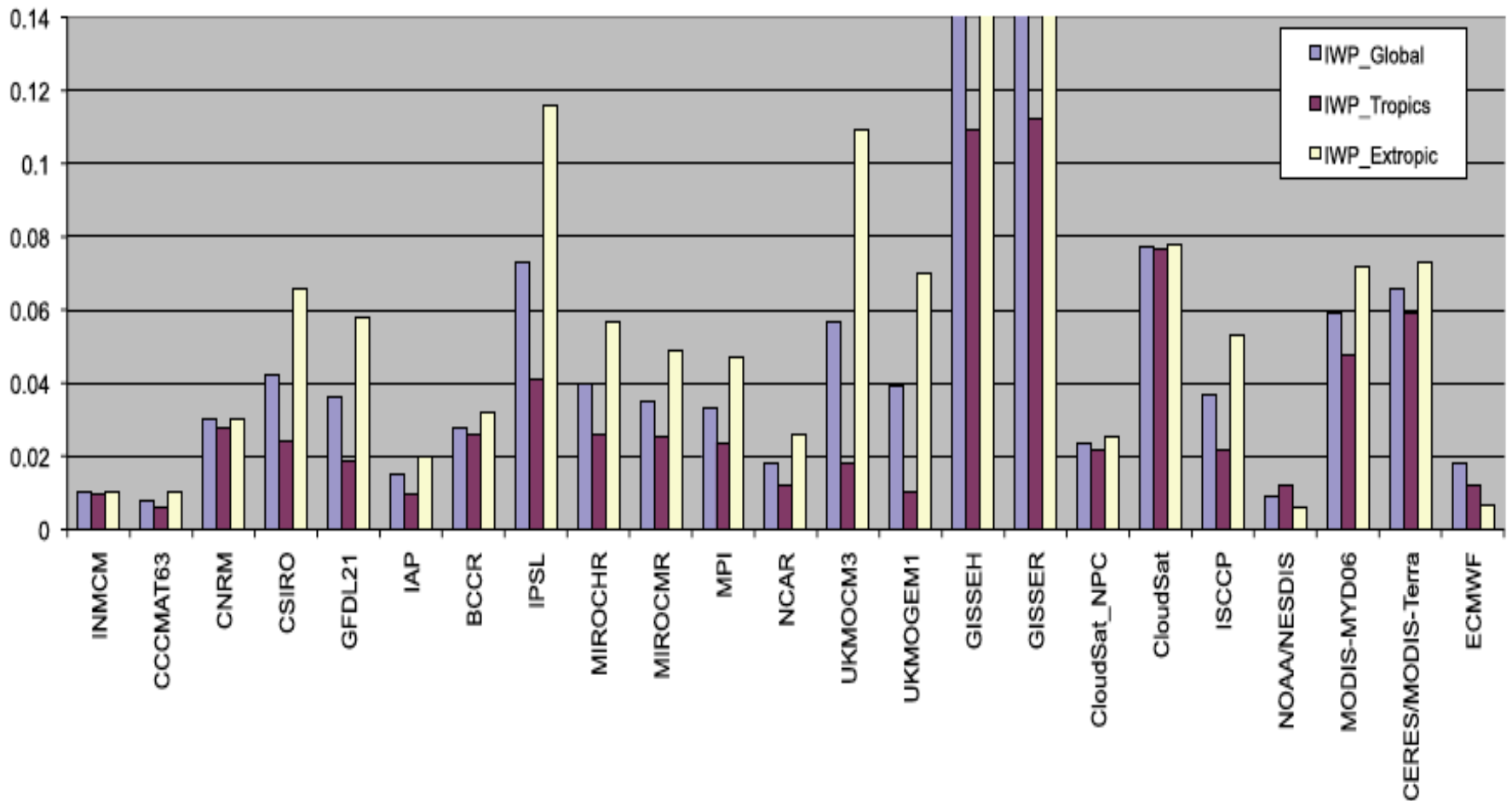
(**THM(Retrieval)/THM(Downstream)** minus **Smooth(Retrieval)/Smooth(Downstream)**)



- Overall regional RMS difference is $\sim 1\%$. However, in some locations regional differences reach 3% .
- Differences tend to be positive in tropics and negative in midlatitudes.

Findings by Loeb et al. (2018): radiative fluxes derived using a consistent ice particle model assumption throughout provide a more robust reference for climate model evaluation compared to existing ice cloud property retrievals.

In other words, the same ice model must be consistently used in **forward remote sensing** implementation (look-up tables) and **downstream radiative forcing assessment**.



Global (blue), tropical (30° N– 30° S; red) and extratropical ($>30^{\circ}$ N,S; yellow) spatial mean values of cloud ice-water path (kg m^{-2}) for 23 GCM simulations (adapted from Waliser et al., 2009). Note that the blue (yellow) bars of GISSER and GISSER that extend above the top of the plot have values of 0.21 and 0.22 (0.34 and 0.36), respectively. Observations are shown in the

Ice Water Path (IWP),

Optical Thickness (tau)

Effective Particle size (D_{eff})

$$\text{IWP} = \text{constant} \cdot \text{tau} \cdot D_{\text{eff}}$$

Summary

- ✦ **Consistency hypothesis:** *radiative fluxes derived using a consistent ice particle model assumption throughout provide a more robust reference for climate model evaluation compared to existing ice cloud property retrievals (Loeb et al. 2017). In other words, the same ice model must be consistently used in forward remote sensing implementations and downstream radiative transfer computations.*

- ✦ **Objectives**
 - 1) Validate the consistency hypothesis by using the MODIS, AIRS, and CALIPSO cloud property products, CERES flux products, and state-of-the-science light scattering and radiative transfer modeling capabilities in conjunction with the ice cloud models used by the respective science teams.
 - 2) Quantify the global scale uncertainties/errors caused by using inconsistent ice cloud models
 - 3) **Develop parameterization schemes** for ice cloud bulk radiative properties that are consistent with the ice models used in NASA's cloud property retrieval products